

## DER HARTREE-FOCK'SCHE VARIATIONS-ANSATZ

$$\hat{H} = \sum_{i,j} K_{ij} \hat{c}_i^\dagger \hat{c}_j + \frac{1}{2} \sum_{i,j,l,m} V_{ijlm} \hat{c}_i^\dagger \hat{c}_j^\dagger \hat{c}_m \hat{c}_l$$

$$K_{ij} = \int dx \varphi_i^*(x) \left[ -\frac{\hbar^2 \nabla^2}{2m} + V_G(\mathbf{r}) \right] \varphi_j(x)$$

$$V_{ijlm} = \int \int dx dx' \varphi_i^*(x) \varphi_j^*(x') V(\mathbf{r} - \mathbf{r}') \varphi_l(x) \varphi_m(x')$$

## VARIATION MIT NEBENBEDINGUNG

$$\delta_\phi \left\{ \langle \phi | \hat{H} | \phi \rangle - \lambda \langle \phi | \phi \rangle \right\} = 0$$

## TRIAL FUNCTION

$$\phi(x_1, x_2, \dots, x_N) = \frac{1}{\sqrt{N!}} \begin{vmatrix} \varphi_{E_1}(x_1) & \varphi_{E_1}(x_2) & \dots & \varphi_{E_1}(x_N) \\ \varphi_{E_2}(x_1) & \varphi_{E_2}(x_2) & \dots & \varphi_{E_2}(x_N) \\ \vdots & \vdots & \ddots & \vdots \\ \varphi_{E_N}(x_1) & \varphi_{E_N}(x_2) & \dots & \varphi_{E_N}(x_N) \end{vmatrix}$$

$$|\phi\rangle \equiv |1_1, 1_2, \dots, 1_N, 0_{N+1}, 0_{N+2}, \dots, 0_\infty\rangle$$

DAS HARTREE-FOCK-POTENTIAL (8.22)

$$V_{HF}^{(i)}(x) = \sum_{j=1, j \neq i}^N \int dx' V(\mathbf{r} - \mathbf{r}') |\varphi_j(x')|^2 - \sum_{j=1, j \neq i}^N \int dx' V(\mathbf{r} - \mathbf{r}') \frac{\varphi_j^*(x') \varphi_i(x') \varphi_j(x)}{\varphi_i(x)}.$$

DER HARTREE-TERM DES POTENTIALS (8.26)

$$V_{Hartree}(\mathbf{r}) = 2 \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} |\varphi_{\mathbf{k}'}(\mathbf{r}')|^2.$$

$$V_{Hartree}(\mathbf{r}) = e^2 \int d^3 r \frac{n_{el}(\mathbf{r})}{|\mathbf{r} - \mathbf{r}'|} \quad \text{mit} \quad n_{el}(\mathbf{r}) = 2 \sum_{\mathbf{k}}^{occ} |\varphi_{\mathbf{k}}(\mathbf{r})|^2$$

DER FOCK- (AUSTAUSCH-) TERM DES POTENTIALS (8.28)

$$V_{Fock}^{\mathbf{k}; \uparrow \uparrow}(\mathbf{r}) = - \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} \frac{\varphi_{\mathbf{k}'}^*(\mathbf{r}') \varphi_{\mathbf{k}}(\mathbf{r}') \varphi_{\mathbf{k}'}(\mathbf{r})}{\varphi_{\mathbf{k}}(\mathbf{r})}.$$

DIE HF-GLEICHUNG IN ORTS-KOORDINATEN (8.30)

$$\begin{aligned}
 & \left[ -\frac{\hbar^2 \nabla^2}{2m} + V_G(\mathbf{r}) + 2 \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} |\varphi_{\mathbf{k}'}(\mathbf{r}')|^2 - \right. \\
 & \left. - \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} \frac{\varphi_{\mathbf{k}'}^*(\mathbf{r}') \varphi_{\mathbf{k}}(\mathbf{r}') \varphi_{\mathbf{k}'}(\mathbf{r})}{\varphi_{\mathbf{k}}(\mathbf{r})} \right] \varphi_{\mathbf{k}}(\mathbf{r}) = \epsilon_{\mathbf{k}}^{HF} \varphi_{\mathbf{k}}(\mathbf{r})
 \end{aligned}$$

DIE GESAMTE HF-SYSTEMENERGIE IN ORTS-KOORDINATEN

(Einsetzen von (8.26) und (8.28) in (8.33))

$$\begin{aligned}
 E_{total} &= 2 \sum_{\mathbf{k}(occ)} \epsilon_{\mathbf{k}}^{HF} - \sum_{\mathbf{k}(occ)} \int d^3 r |\varphi_{\mathbf{k}}(\mathbf{r})|^2 \\
 &\times \left[ 2 \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} |\varphi_{\mathbf{k}'}(\mathbf{r}')|^2 - \right. \\
 &\left. - \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} \frac{\varphi_{\mathbf{k}'}^*(\mathbf{r}') \varphi_{\mathbf{k}}(\mathbf{r}') \varphi_{\mathbf{k}'}(\mathbf{r})}{\varphi_{\mathbf{k}}(\mathbf{r})} \right] + \\
 &+ \frac{1}{2} \int d^3 r d^3 r' \frac{\rho_G(\mathbf{r}) \rho_G(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}
 \end{aligned}$$

## DIE HARTREE-FOCK-THEORIE FÜR FREIE ELEKTRONEN

$$\rho_G(\mathbf{r}) \approx \rho_G = \frac{N e}{\Omega} \rightarrow V_G(\mathbf{r}) = -\frac{N e^2}{\Omega} \int_{(\Omega)} \frac{d^3 r'}{|\mathbf{r} - \mathbf{r}'|}.$$

## DIE HARTREE-FOCK-GLEICHUNG FÜR JELLIUM (8.36)

$$\left\{ -\frac{\hbar^2 \nabla^2}{2m} - \frac{N e^2}{\Omega} \int_{(\Omega)} \frac{d^3 r'}{|\mathbf{r} - \mathbf{r}'|} + 2 \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} |\varphi_{\mathbf{k}'}(\mathbf{r}')|^2 - \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} \frac{\varphi_{\mathbf{k}'}^*(\mathbf{r}') \varphi_{\mathbf{k}}(\mathbf{r}') \varphi_{\mathbf{k}'}(\mathbf{r})}{\varphi_{\mathbf{k}}(\mathbf{r})} \right\} \varphi_{\mathbf{k}}(\mathbf{r}) = \varepsilon_{\mathbf{k}}^{HF} \varphi_{\mathbf{k}}(\mathbf{r}).$$

EIN EBENE-WELLEN-ANSATZ:

$$\varphi_{\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{\Omega}} e^{i\mathbf{k} \cdot \mathbf{r}}$$

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$$\left\{ -\frac{\hbar^2 \nabla^2}{2m} - \frac{N e^2}{\Omega} \int_{(\Omega)} \frac{d^3 r'}{|\mathbf{r} - \mathbf{r}'|} + \frac{2}{\Omega} \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} - \frac{1}{\Omega} \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} e^{i(\mathbf{k}' - \mathbf{k})\cdot(\mathbf{r} - \mathbf{r}')} \right\} e^{i\mathbf{k}\cdot\mathbf{r}} = \varepsilon_{\mathbf{k}}^{HF} e^{i\mathbf{k}\cdot\mathbf{r}}.$$

$$\left\{ -\frac{\hbar^2 \nabla^2}{2m} - \frac{1}{\Omega} \sum_{\mathbf{k}'(occ)} \int d^3 r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} e^{i(\mathbf{k}' - \mathbf{k})\cdot(\mathbf{r} - \mathbf{r}')} \right\} e^{i\mathbf{k}\cdot\mathbf{r}} = \varepsilon_{\mathbf{k}}^{HF} e^{i\mathbf{k}\cdot\mathbf{r}}.$$

$$\left\{ \frac{\hbar^2 k^2}{2m} - \underbrace{\frac{1}{\Omega} \sum_{\mathbf{k}'(occ)} \int d^3 z \frac{e^2}{z} e^{-i(\mathbf{k}' - \mathbf{k})\cdot\mathbf{z}}}_{\text{Jellium-Austauschenergie}} \right\} e^{i\mathbf{k}\cdot\mathbf{r}} = \varepsilon_{\mathbf{k}}^{HF} e^{i\mathbf{k}\cdot\mathbf{r}}.$$

$$\epsilon_{\mathbf{k}}^x = -\frac{e^2 k_F}{2\pi} \left[ 2 + \frac{(k_F^2 - k^2)}{k \cdot k_F} \cdot \ln \left| \frac{k + k_F}{k - k_F} \right| \right] \quad (8.43)$$

## HF-DISPERSION FÜR JELLIUM:

$$\epsilon_{\mathbf{k}}^{HF} = \frac{\hbar^2 k^2}{2m} + \epsilon_{\mathbf{k}}^x$$

bzw.

$$\epsilon_{|\mathbf{k}|}^{HF} = \frac{\hbar^2 k^2}{2m} - \frac{e^2 k_F}{2\pi} \left[ 2 + \frac{(k_F^2 - k^2)}{k \cdot k_F} \cdot \ln \left| \frac{k + k_F}{k - k_F} \right| \right] \quad (8.44)$$

## DIE GESAMTE HF-SYSTEMENERGIE IN ORTS-KOORDINATEN

(Einsetzen von (8.26) und (8.28) in (8.33))

$$\begin{aligned}
 E_{total} = & 2 \sum_{\mathbf{k}(occ)} \epsilon_{\mathbf{k}}^{HF} - \sum_{\mathbf{k}(occ)} \int d^3r |\varphi_{\mathbf{k}}(\mathbf{r})|^2 \\
 & \times \left[ 2 \sum_{\mathbf{k}'(occ)} \int d^3r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} |\varphi_{\mathbf{k}'}(\mathbf{r}')|^2 - \right. \\
 & \left. - \sum_{\mathbf{k}'(occ)} \int d^3r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} \frac{\varphi_{\mathbf{k}'}^*(\mathbf{r}') \varphi_{\mathbf{k}}(\mathbf{r}') \varphi_{\mathbf{k}'}(\mathbf{r})}{\varphi_{\mathbf{k}}(\mathbf{r})} \right] + \\
 & + \frac{1}{2} \int d^3r d^3r' \frac{\rho_G(\mathbf{r}) \rho_G(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}
 \end{aligned}$$

mit

$$\rho_G(\mathbf{r}) \approx \rho_G = \frac{N e}{\Omega} \quad \text{und} \quad \varphi_{\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{\Omega}} e^{i\mathbf{k} \cdot \mathbf{r}}$$

$$\begin{aligned}
 E_{total} = & 2 \sum_{\mathbf{k}(occ)} \epsilon_{\mathbf{k}}^{HF} - \frac{1}{\Omega^2} \sum_{\mathbf{k}(occ)} \int d^3r \, 2 \sum_{\mathbf{k}'(occ)} \int d^3r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} + \\
 & \frac{1}{\Omega^2} \sum_{\mathbf{k}(occ)} \int d^3r \sum_{\mathbf{k}'(occ)} \int d^3r' \frac{e^2}{|\mathbf{r} - \mathbf{r}'|} e^{i(\mathbf{k}' - \mathbf{k}) \cdot (\mathbf{r} - \mathbf{r}')} + \\
 & + \frac{1}{2} \left( \frac{N e}{\Omega} \right)^2 \iint \frac{d^3r d^3r'}{|\mathbf{r} - \mathbf{r}'|}.
 \end{aligned}$$

Wieder heben sich im Fall Jellium die **roten** Terme weg, und es ergibt sich

$$E_{total} = 2 \sum_{\mathbf{k}(occ)} \varepsilon_{\mathbf{k}}^{HF} - \frac{1}{\Omega} \sum_{\mathbf{k}(occ)} \int d^3r \left\{ -\frac{1}{\Omega} \sum_{\mathbf{k}'(occ)} \int d^3z \frac{e^2}{z} e^{i(\mathbf{k}'-\mathbf{k})\cdot\mathbf{z}} \right\}$$

bzw.

$$E_{total} = 2 \sum_{\mathbf{k}(occ)} \varepsilon_{\mathbf{k}}^{HF} - \sum_{\mathbf{k}(occ)} \varepsilon_{\mathbf{k}}^x.$$

Unter Berücksichtigung von (8.39) ergeben sich die Gleichungen

$$E_{total} = 2 \sum_{\mathbf{k}(occ)} \frac{\hbar^2 k^2}{2m} + 2 \sum_{\mathbf{k}(occ)} \varepsilon_{\mathbf{k}}^x - \sum_{\mathbf{k}(occ)} \varepsilon_{\mathbf{k}}^x$$

bzw.

$$E_{total} = \underbrace{2 \sum_{\mathbf{k}(occ)} \frac{\hbar^2 k^2}{2m}}_{\text{kinetische Energie}} + \underbrace{\sum_{\mathbf{k}(occ)} \varepsilon_{\mathbf{k}}^x}_{\text{Austauschenergie}} \quad (8.46)$$

Auswertung von (8.46) führt zur Gleichung (8.49):

$$\frac{E_{total}}{N} = \left( \frac{\hbar^2}{2m} \right) \left( \frac{9\pi}{4} \right)^{2/3} \frac{3}{5r_s^2} - \left( \frac{9\pi}{4} \right)^{1/3} \frac{3e^2}{4\pi r_s}$$