


# Quantum Impurity Models in and out of equilibrium studied by means of Variational Cluster Perturbation Theory

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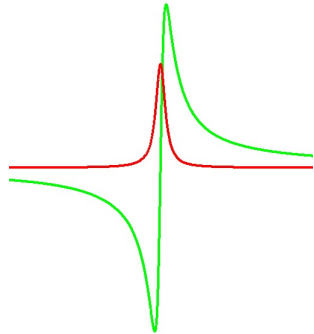
## Generic lattice model

$$\begin{aligned}\hat{\mathcal{H}} &= \sum_{ij} t_{ij} c_i^\dagger c_j + \sum_{ijkl} U_{ijkl} c_i^\dagger c_j^\dagger c_k c_l \\ &= \hat{\mathcal{H}}_I(t) + \hat{\mathcal{H}}_{II}(U)\end{aligned}$$

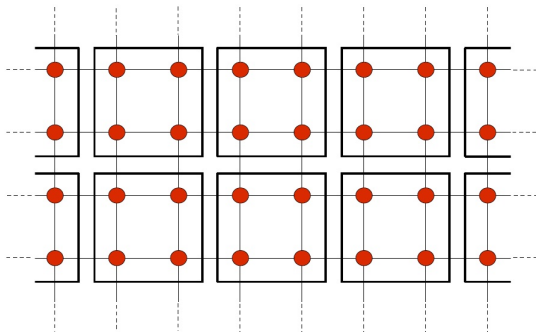
To apply a perturbative method one has to expand in some parameter. Usually the **hopping (strong coupling)** or the **interaction (weak coupling)** are considered.

# Atomic limit

$$G(z; t, U) = \left\langle \frac{1}{z - \hat{\mathcal{H}}(t, U)} \right\rangle$$



- strong coupling perturbation theory (in hopping  $t$ )
- consider  $G(t = 0)$  as starting point



Extrapolate cluster to thermodynamic limit:

- Cluster Perturbation Theory (**CPT**)
- **Variational Cluster Approach (VCA)**
- Cluster/Cellular Dynamical Mean-Field Theory (**CDMFT**)
- Dynamical Cluster Approximation (**DCA**)

# Cluster Perturbation Theory<sup>a b</sup>

<sup>a</sup>C. Gros and R. Valentí, Phys. Rev. B 48, 418 (1993) (I)

<sup>b</sup>D. Sénéchal, D. Perez, and M. Pioro-Ladrière, Phys. Rev. Lett. 84, 522 (2000)

By means of strong coupling perturbation theory it can be shown that the first order result for the lattice Green function  $G$  is

$$G^{-1}(\omega, \mathbf{k}) = G_{\text{cluster}}^{-1}(\omega) - T(\mathbf{k})$$

- $G_{\text{cluster}}$  = exact Green's function of the **cluster**
- $T$  = **inter-cluster** off diagonal one particle terms (i.e. hopping)

# Cluster Perturbation Theory - Limits

CPT is exact for

- $t \rightarrow 0$ ,
- $U \rightarrow 0$ ,
- $L \rightarrow \infty$ .

CPT captures **short-range correlations** exactly, long-range correlations are neglected.

# Variational Cluster Approach<sup>a</sup>

<sup>a</sup>M. Potthoff, M. Aichhorn, and C. Dahnken, *Phys. Rev. Lett.* **91**, 206402 (2003)

- VCA = variational extension to CPT - rigorously developed within the **Self-Energy Functional Approach (SFA)**<sup>ab</sup>,
- does **not** implement a variational principle in the sense of a Rayleigh-Ritz variational principle,
- is applicable to **broken-symmetry/ordered** phases.

<sup>a</sup>M. Potthoff, *Eur. Phys. J. B* **32**, 429 (2003)

<sup>b</sup>M. Potthoff, *Eur. Phys. J. B* **36**, 335 (2003)



# Luttinger-Ward functional $\Phi[G]$

- $\Phi[\Omega]$  = sum of all **two-particle irreducible** diagrams

$$\Phi = \begin{array}{c} \circ \\ | \\ \circ \end{array} + \text{---} \circ \text{---} + \begin{array}{c} \circ \\ \text{---} \\ \circ \end{array} + \dots$$

- The **functional derivative** of  $\Phi[G]$  is the self-energy:

$$\frac{\delta\Phi[G]}{\delta G} = \Sigma$$

- It is a **universal functional** of  $G$ .

## Legendre Transform of $\Phi[G]$

It can be shown that

$$\frac{\delta\Phi[G]}{\delta G} = \Sigma,$$

is **locally invertible**.

**Legendre Transform** of the Luttinger-Ward functional:

$$F[\Sigma] = \Phi[\Sigma] - \text{Tr} \{ \Sigma G \} .$$

It can be shown that the generalized **grand potential functional** is given by:

$$\Omega[\Sigma] = F[\Sigma] - \text{Tr} \{ (-G_0^{-1} + \Sigma) \} .$$

# Grand Potential Functional $\Omega[\Sigma]$

Dyson's equation is recovered at the stationary point of the grand potential functional  $\Omega[\Sigma]$

$$\frac{\delta F[\Sigma]}{\delta \Sigma} = -G$$
$$\frac{\delta \Omega[\Sigma]}{\delta \Sigma} = -G + (G_0^{-1} - \Sigma)^{-1} \stackrel{!}{=} 0.$$

## VCA Reference System (I)

Since systems which share the same interaction part  $\hat{\mathcal{H}}_{II}(U)$  have the same  $\Phi[G]$  (or  $F[\Sigma]$ ) we construct a **reference system**:

$$\hat{\mathcal{H}}' = \hat{\mathcal{H}}_I(t') + \hat{\mathcal{H}}_{II}(U)$$

- defined on the **same lattice**,
- having the **same interaction** as the original system,
- but may have entirely different **single-particle operators / parameters**.

## VCA Reference System (II)

The reference system  $\hat{\mathcal{H}}'$  may be used to **eliminate the Luttinger-Ward functional**: (This is still exact!)

$$\Omega[\Sigma] = F[\Sigma] - \text{Tr} \{ \ln (-G_0^{-1} + \Sigma) \}$$

$$\Omega'[\Sigma] = F[\Sigma] - \text{Tr} \{ \ln (-G_0'^{-1} + \Sigma) \}$$


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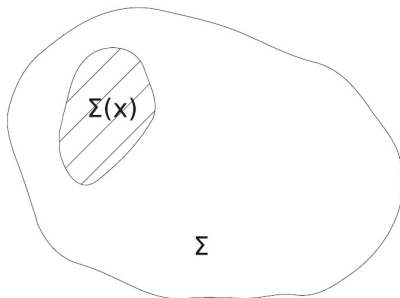
$$- \sum$$

$$\Omega[\Sigma] = \Omega'[\Sigma] + \text{Tr} \{ \ln (-G_0'^{-1} + \Sigma) \} - \text{Tr} \{ \ln (-G_0^{-1} + \Sigma) \}$$

$$= \Omega'[\Sigma] + \text{Tr} \{ \ln (-G'[\Sigma]) \} - \text{Tr} \{ \ln (-G[\Sigma]) \}$$

# Evaluation of the grand potential

The self-energy  $\Sigma(x)$  is given by the self-energy of the reference system  $\Sigma(x')$ , where  $x$  denotes the single particle parameters, **restricting the space of available self-energies**.



- **Grand potential**

$$\Omega(x') = \Omega'(x') + \text{Tr} \{ \ln (-G(x')) \} - \text{Tr} \{ \ln (-G'(x')) \} .$$

- **Stationarity condition:**

$$\nabla_{x'} \Omega(x') \stackrel{!}{=} 0 .$$

- The **Green's function** of the physical system  $G$  is obtained in CPT/VCA by Dyson's equation

$$G^{-1} = G'^{-1} - T .$$

The matrix  $T = G_0'^{-1} - G_0^{-1}$  contains all single particle terms not included in the reference system (as in CPT) as well as the deviation, introduced by VCA,  $\Delta x \equiv x' - x$  of the single-particle parameters of the reference system  $x'$  with respect to the ones of the original system  $x$ .



$$\text{Seminar Appeal} = \frac{\text{Relevance} \times \text{Food}}{(\text{Distance})^2}$$

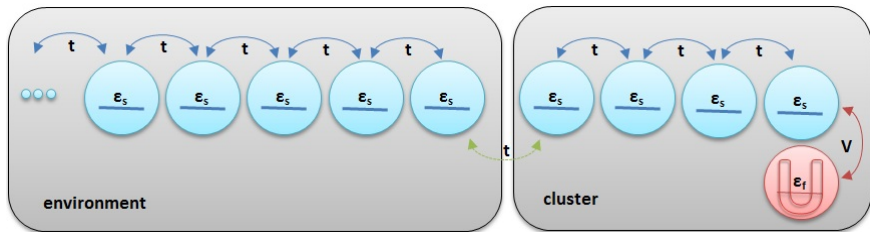


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# Single Impurity Anderson Model - Cluster decomposition



# Grand potential for infinite reference systems

The reference system here consists of **two parts**:

- a **cluster part** and
- an **infinite environment**.

It is possible to find a **regular expression for the grand potential** (at zero temperature):

$$\Omega - \Omega'_{0,\text{env}} = \Omega'_{0,\text{cluster}} + \text{tr} \{T\} - \frac{1}{\pi} \sum_{\sigma} \int_0^{\infty} d\omega \ln \left| \det \left( \mathbb{1}_{cc} - \tilde{\Sigma}_{cc}(i\omega) G'_{cc}(i\omega) \right) \right| .$$

# Numerical implementation

The Green's function of the reference system are obtained by

- **Lanczos/Band-Lanczos method** for the cluster part,
- **analytically** for the environment part.

# Results in equilibrium

## Results

- **Spectral properties** (comparison to NRG, FRG<sup>a</sup>, DMRG<sup>b</sup>, CTQMC),
- **Impurity density of states and occupation** (Friedel sum rule) (comparison to Hubbard-X operator technique results<sup>c</sup>, "Crossover" diagram (comparison to mean field results<sup>d</sup>))
- **Low energy properties**, Kondo temperature, effective mass, static spin susceptibility

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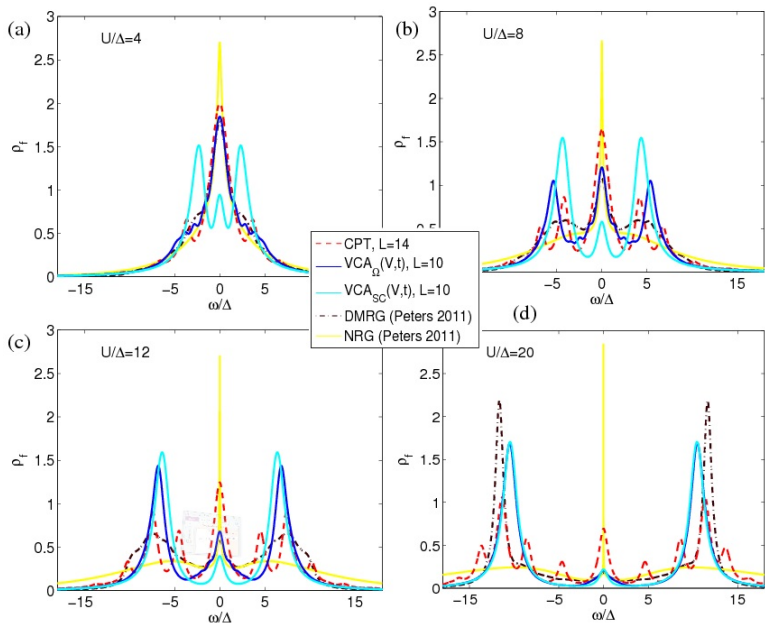
<sup>a</sup>C. Karrasch, R. Hedden, R. Peters, T. Pruschke, K. Schönhammer, and V. Meden, *J. Phys.: Condensed Matter* 20, 345205 (2008)

<sup>b</sup>R. Peters, 1103.5837 (2011)

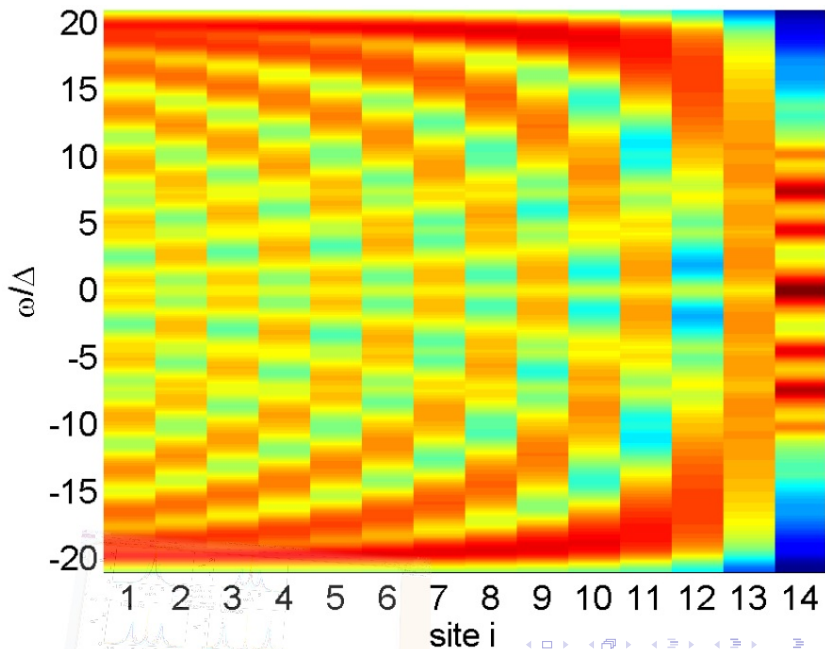
<sup>c</sup>T. Lobo, M. S. Figueira, and M. E. Foglio, *Nanotechnology* 21, 274007 (2010)

<sup>d</sup>P. Coleman, *AIP Conference Proceedings* 629, 79 (2002)

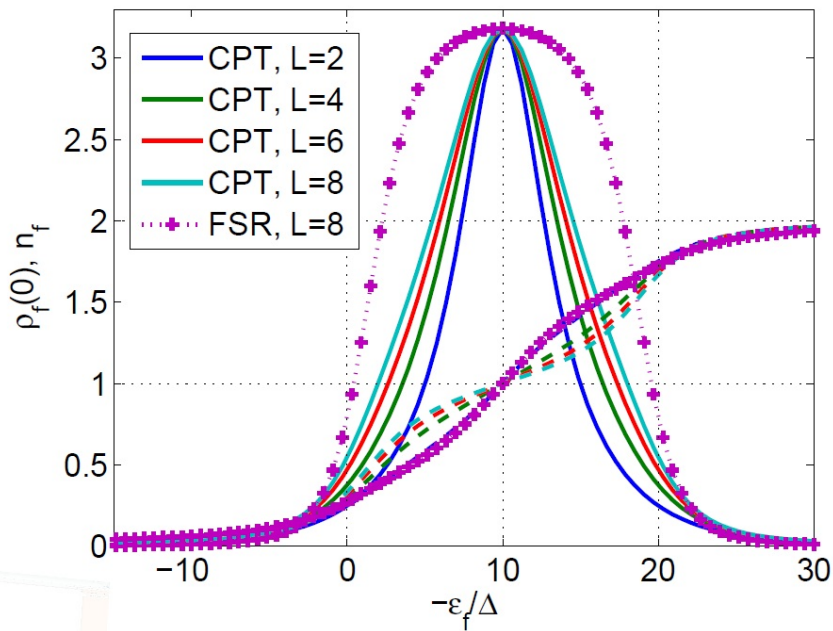
# Spectral properties



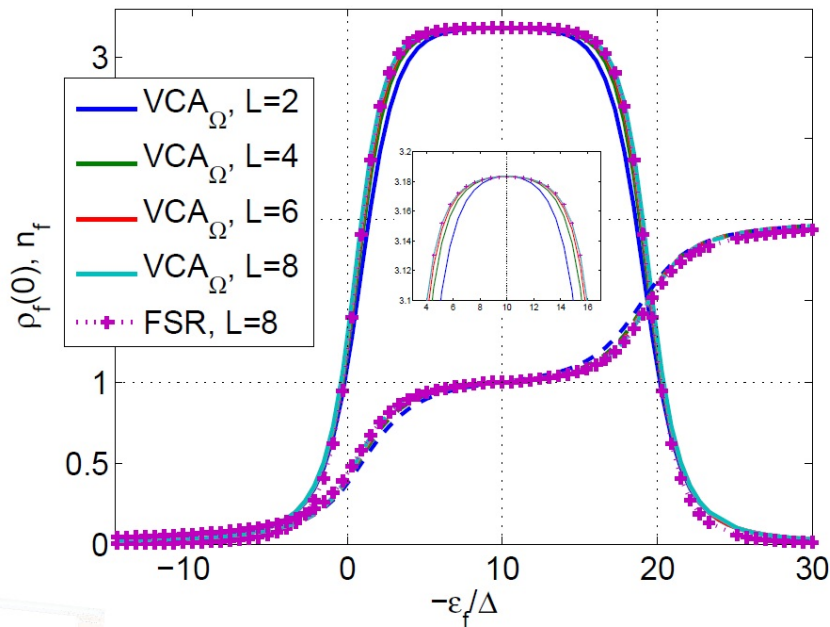
# Spectral properties (spatial)



# Impurity Occupation

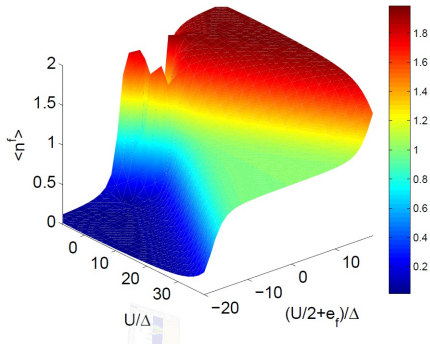
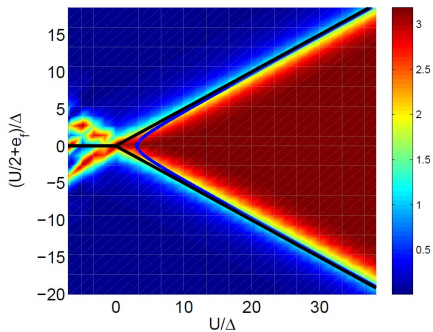


# Impurity Occupation





# "Crossover" diagram



# Kondo Temperature

- Since the height of the Kondo resonance is fixed by the Friedel sum rule, the **spectral weight** (area) of Kondo peak and its **FWHM** are **proportional** to  $T_K$ .
- The Kondo temperature (symmetric SIAM) is given by **Bethe Ansatz**<sup>2</sup>

$$T_K = \sqrt{\frac{\Delta U}{2}} e^{-\gamma \frac{\pi}{8\Delta} U}, \quad \gamma = 1.$$

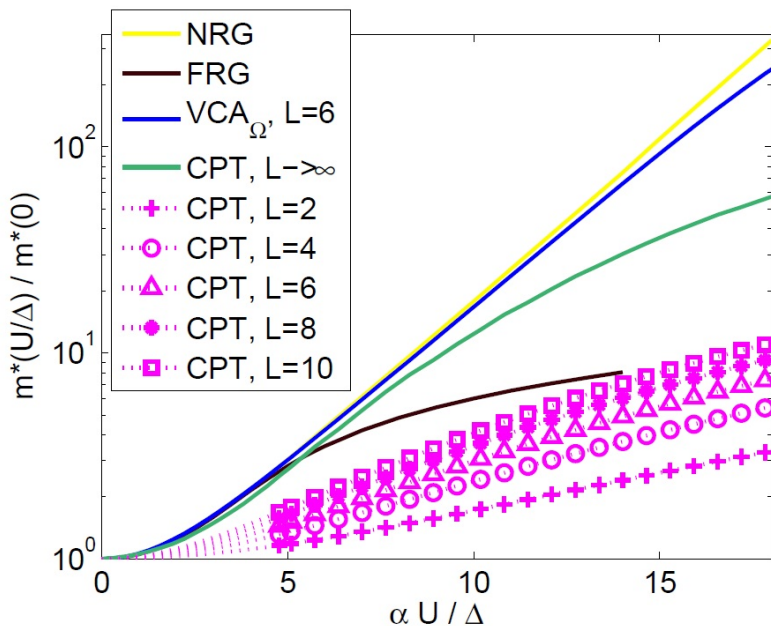
- An analytic calculation for a two-site reference system yields for VCA<sub>Ω</sub>

$$\gamma = 0.6511.$$

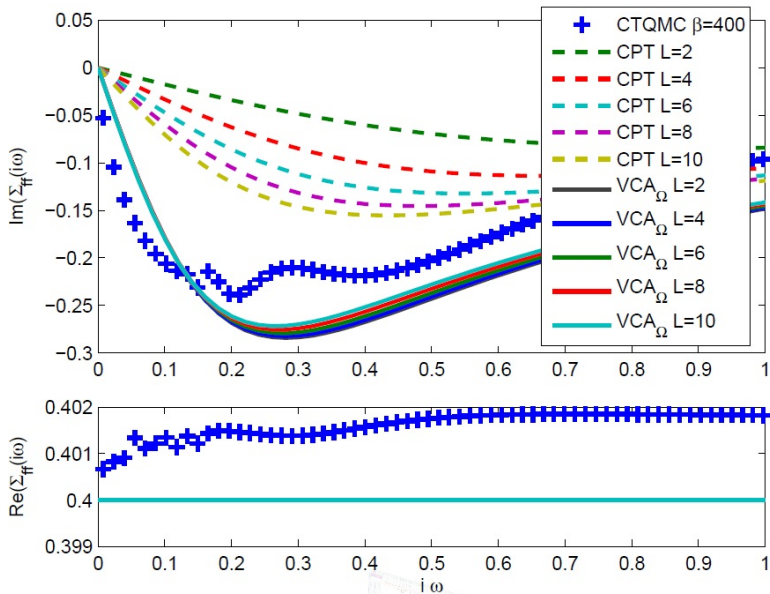
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<sup>2</sup>A. C. Hewson, *The Kondo Problem to Heavy Fermions* (Cambridge University Press, 1997)

# Effective Mass



# Comparison to CTQMC



# Conclusion

- **VCA** » CPT
- **Kondo peak** + exponential scale in  $U$
- **Hubbard bands** (position + width)
- **all parameter regions** (pinning of Kondo resonance)
- $\Sigma$  exact for high Matsubara frequency
- extension to **many orbitals, arbitrary dimensions, non-equilibrium** feasible
- **fast**

# Nonequilibrium extension of VCA (I)

- Initial state: **three decoupled systems in equilibrium**  
left lead - cluster - right lead.
- At some time  $t_0$  the **coupling is switched on**.
- We are interested in the long time **steady-state properties**.

# Nonequilibrium extension of VCA (II)

- **Keldysh - formalism** to obtain steady-state properties.
- **VCA reformulated** in terms of self-consistently determined variational parameters where the self-consistency conditions are static expectation values (for example):

$$\begin{aligned} \langle \hat{n}_\sigma^f \rangle_{\text{cluster}, \epsilon'_f, \epsilon'_s} &\stackrel{!}{=} \langle \hat{n}_\sigma^f \rangle_{\text{CPT}, \epsilon_f, \epsilon_s, \epsilon'_f, \epsilon'_s} \\ \sum_i^{L-1} \langle \hat{n}_\sigma^i \rangle_{\text{cluster}, \epsilon'_f, \epsilon'_s} &\stackrel{!}{=} \sum_i^{L-1} \langle \hat{n}_\sigma^i \rangle_{\text{CPT}, \epsilon_f, \epsilon_s, \epsilon'_f, \epsilon'_s} \end{aligned} \cdot$$

- Green's functions calculated in **Keldysh space on the real energy axis**:

$$\tilde{G}(\omega) = \begin{pmatrix} G^{\text{ret}}(\omega) & G^{\text{keld}}(\omega, \mu) \\ 0 & G^{\text{adv}}(\omega) \end{pmatrix}$$

- The initial  $G^{\text{keld}}(\omega, \mu)$  of the decoupled system is given by

$$G^{\text{keld}}(\omega, \mu) = \left( G^{\text{ret}}(\omega) - G^{\text{adv}}(\omega) \right) (1 - 2p_{\text{FD}}(\omega, \mu, \beta))$$

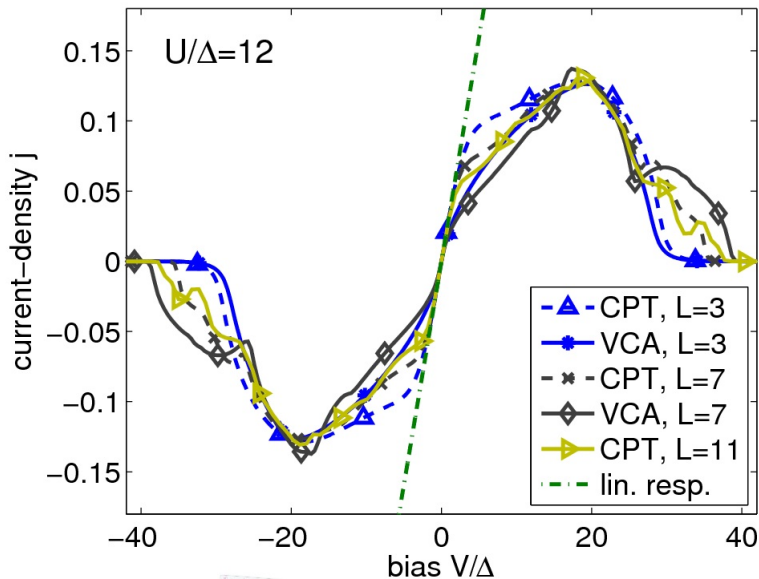
# Current for a single impurity orbital

The expression for the current between sites  $i$  and  $j$  is given by:

$$\begin{aligned} I_{ij} &= t \operatorname{Re} \left( G_{ij}^{\text{keld}}(t=0) \right) \\ &= \frac{t}{2} \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} \operatorname{Re} \left( G_{ij}^{\text{keld}}(\omega) - G_{ji}^{\text{keld}}(\omega) \right) \end{aligned}$$



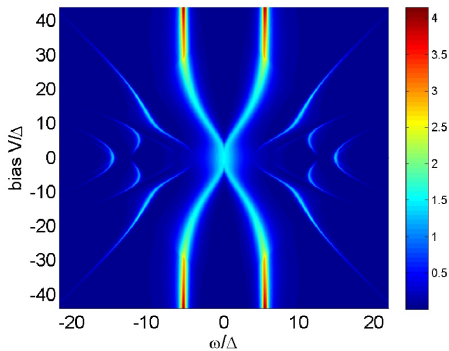
# Current for a single impurity orbital



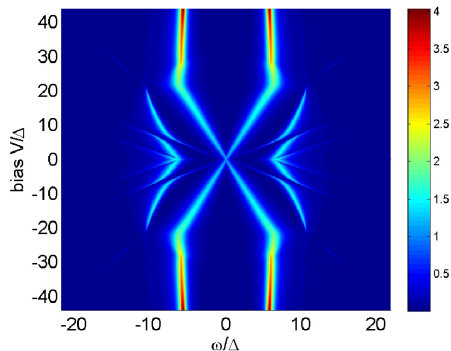
# Non-equilibrium density of states

$$\frac{U}{\Delta} = 12$$

CPT



VCA



# Thank You!

Thank you for your attention!

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