

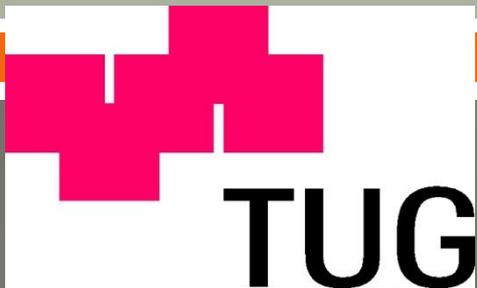


Strongly correlated quantum systems out of equilibrium

a non-equilibrium variational cluster approach



Martin Nuss, supervised by Wolfgang von der Linden & Enrico Arrigoni



Ljubljana, 11.12.2012





Christoph Heil

Prof. Wolfgang von der Linden



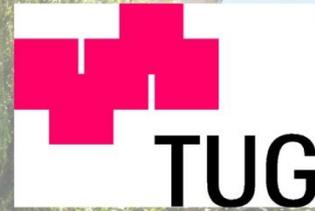
Michael Knap

Prof. Enrico Arrigoni



Martin Nuss

non-equilibrium group at



Prof. Hans Gerd Evertz



Martin Ganahl



Agenda

1

non-equilibrium phenomena - Quantum Dots

2

non-equilibrium Variational Cluster Approach

3

Matrix Product State time evolution

4

Results: steady – state

5

Results: time evolution

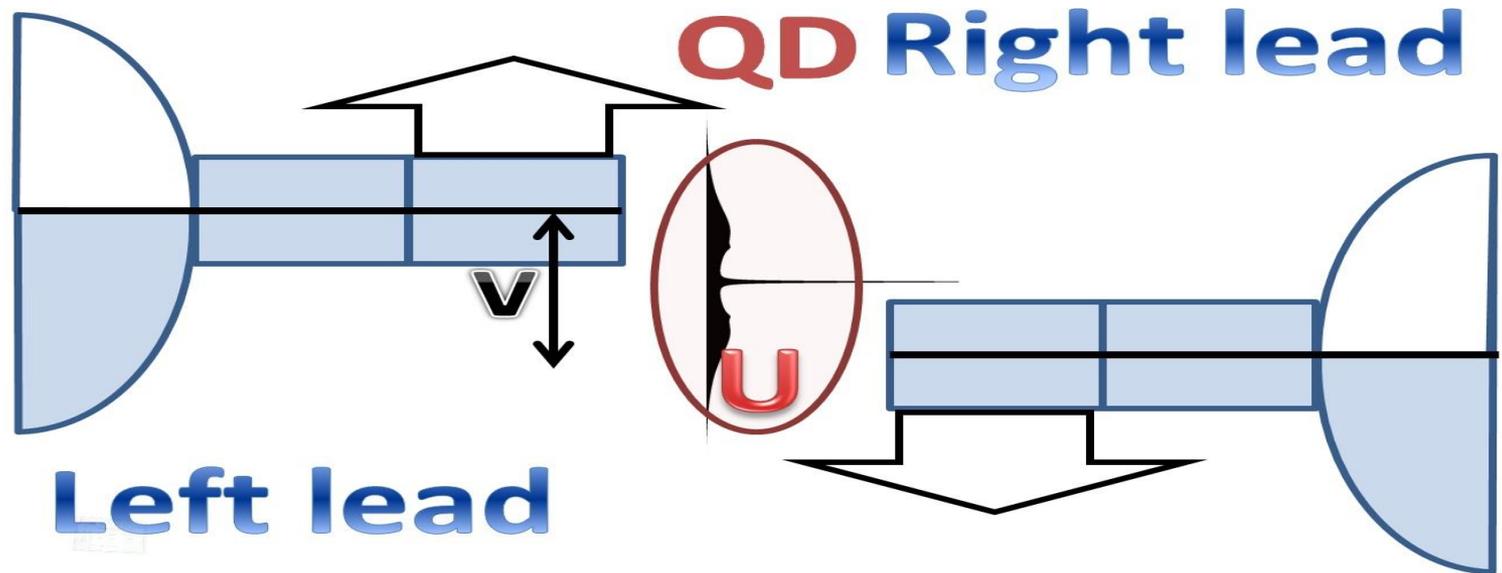
Non-equilibrium phenomena — Quantum Dots

∞ 1 ∞

Transport through a strongly correlated object

Single Impurity Anderson Model

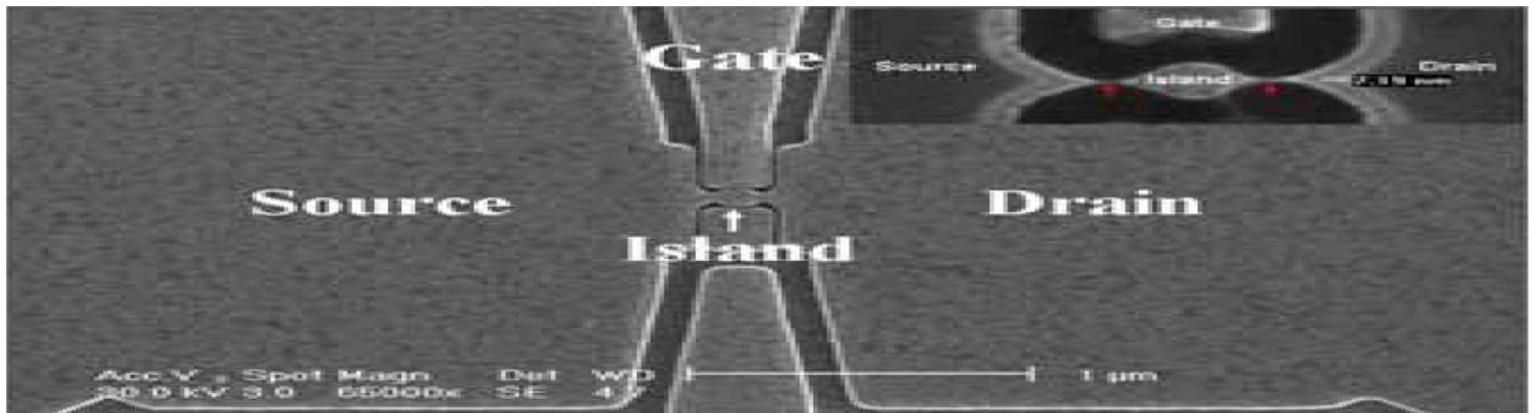
P. W. Anderson (1961)



+

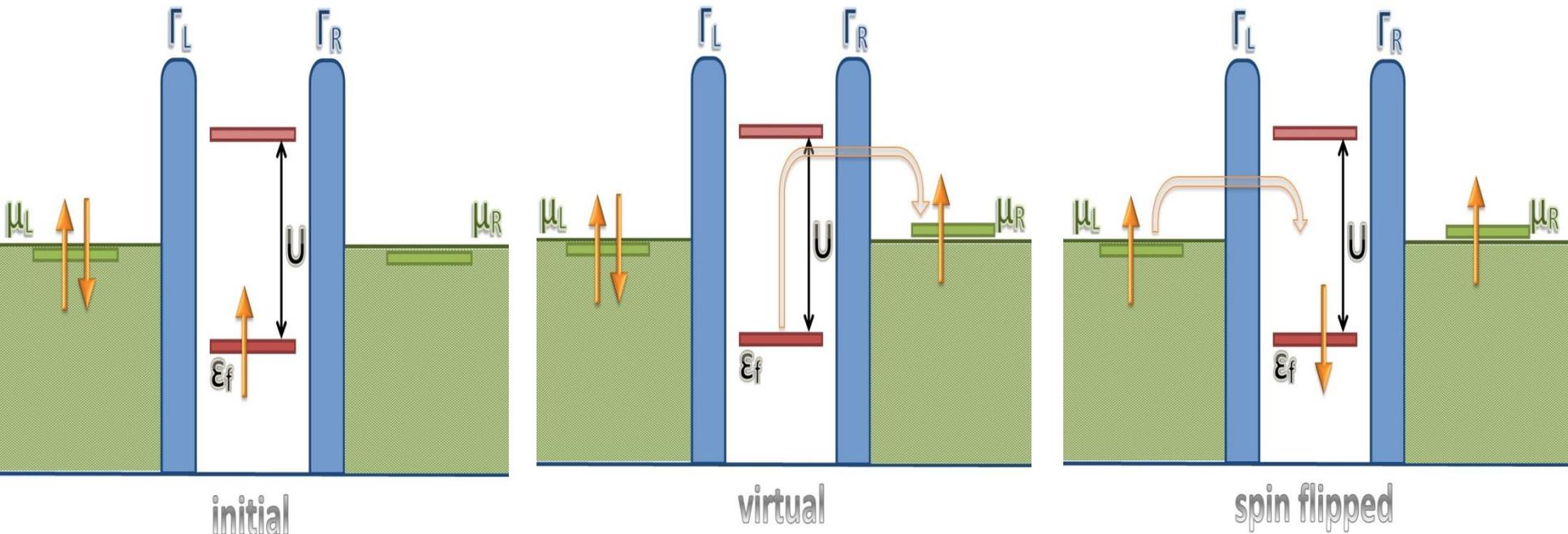
voltage bias

11.12.2012



C. S. Wu, Y. Makiuchi, C. D. Chen, ed. M. Wang (2010)

Transport through a strongly correlated object



Y. Meir, N.S. Wingreen (1992)

$$\mathbf{J} = \frac{ie}{2h} \int d\epsilon (\text{tr}\{[f_L(\epsilon)\Gamma^L - f_R(\epsilon)\Gamma^R](\mathbf{G}^r - \mathbf{G}^a)\} + \text{tr}\{(\Gamma^L - \Gamma^R)\mathbf{G}^<\})$$

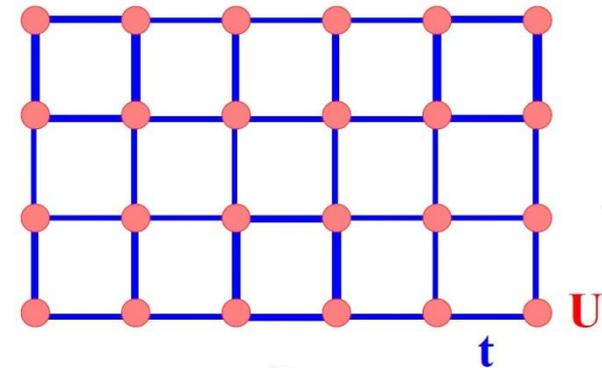
need:

single particle Green's function G in Keldysh space

Non-equilibrium Variational Cluster Approach

∞ 2 ∞

Many-Body cluster methods



given $\hat{\mathcal{H}}$

ask for \mathbf{G}

in general **unsolvable**



Strategy ?

1)

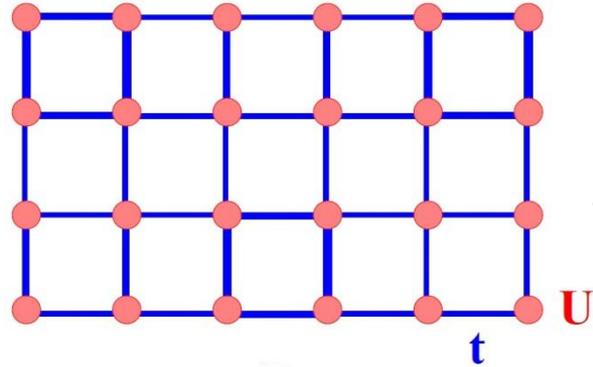
2)

3)

4)



Many-Body cluster methods



given $\hat{\mathcal{H}}$

ask for \mathbf{G}

in general **unsolvable**



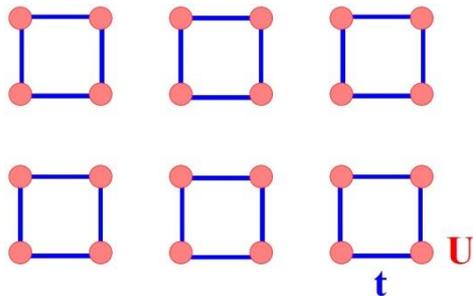
Strategy ?

1) **CUT**

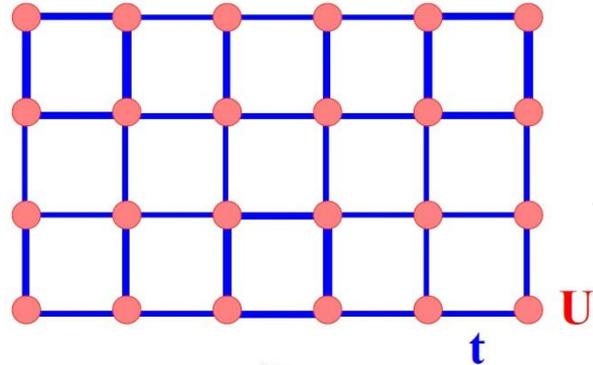
2)

3)

4)



Many-Body cluster methods



given $\hat{\mathcal{H}}$

ask for \mathbf{G}

in general **unsolvable**



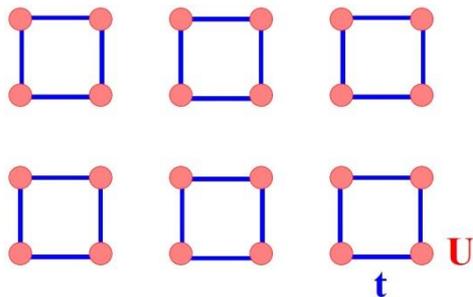
Strategy ?

1) **CUT**

2) **SOLVE**

3)

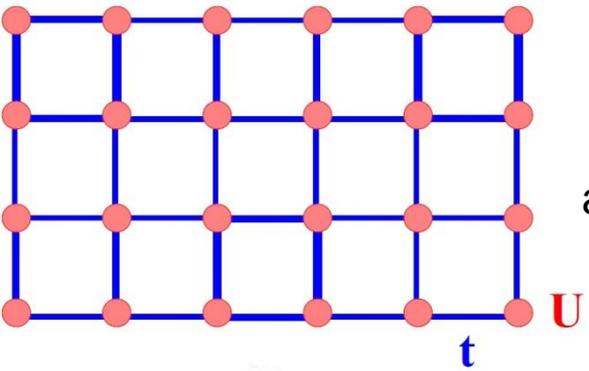
4)



$\mathbf{G}_{\text{cluster}}$
=
exactly solvable



Cluster Perturbation Theory (CPT)



given $\hat{\mathcal{H}}$
ask for \mathbf{G}

• **Cluster Perturbation Theory (CPT)**

- C. Gros, R. Valentí (1993)
- D. Sénéchal, D. Perez, M. Pioro-Ladrière (2000)

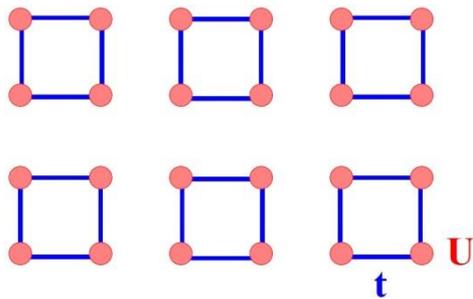
first order strong coupling perturbation theory

$$\Sigma = \sum_{\text{cluster}}$$

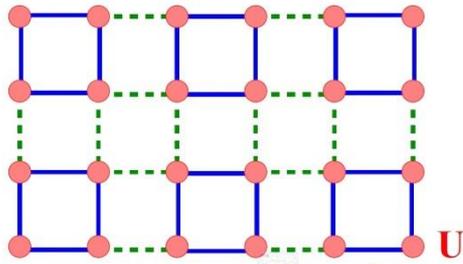
1) **CUT**

2) **SOLVE**

3) **GLUE**



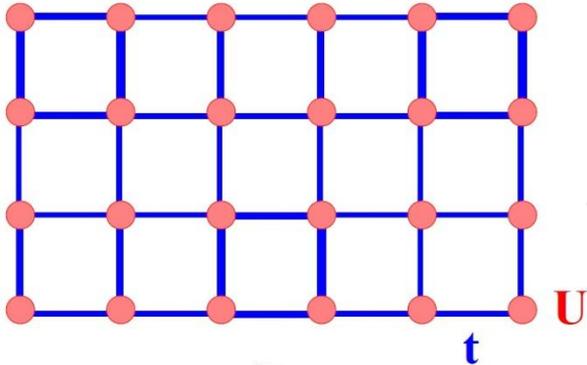
$\mathbf{G}_{\text{cluster}}$
=
exactly solvable



$$\mathbf{G} = \mathbf{G}_{\text{cluster}}^{-1} - \mathbf{T}$$



Variational Cluster Approach (VCA)



given $\hat{\mathcal{H}}$
ask for \mathbf{G}

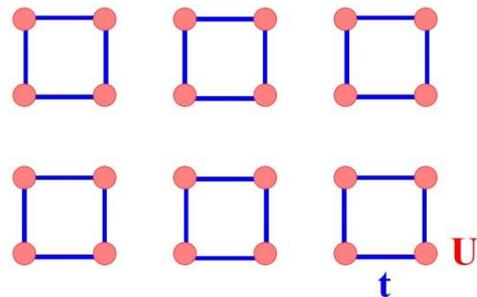
- **Cluster Perturbation Theory (CPT)**
C. Gros, R. Valenti (1993)
D. Sénéchal, D. Perez, M. Pioro-Ladrière (2000)
- **Variational Cluster Approach (VCA)**
M. Potthoff, M. Aichhorn, C. Dahnen (2003)
- Cellular Dynamical Mean-Field Theory (CDMFT)
- Dynamical Cluster Approximation (DCA)

1) CUT

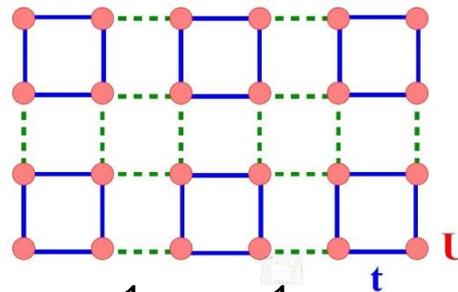
2) SOLVE

3) GLUE

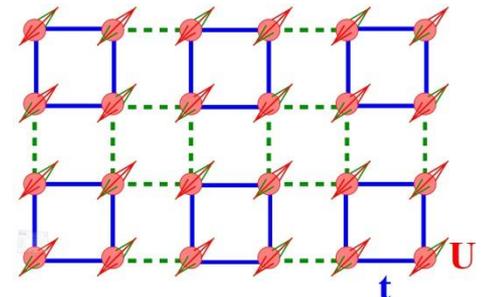
4) ADD FIELDS



$\mathbf{G}_{\text{cluster}}$
=
exactly solvable



$$\mathbf{G} = \mathbf{G}_{\text{cluster}}^{-1} - \mathbf{T}$$



$$\Sigma = \sum_{\text{cluster}} \Sigma(\mathbf{x})$$

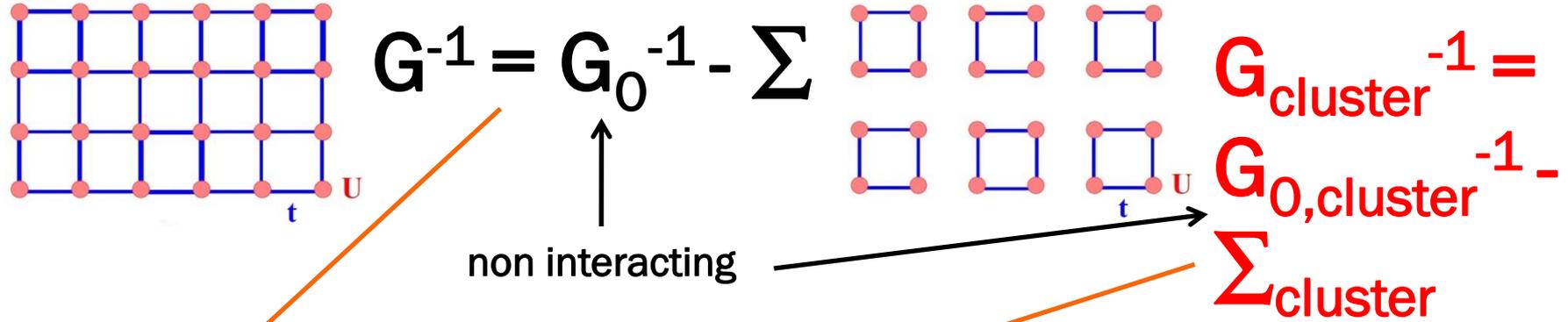
M. Potthoff (2003)

+ Variational principle to fix \mathbf{x} : Self-energy Functional Approach 12

Equilibrium Theory

2 - A 3

Motivation for CPT



$G^{-1} = G_0^{-1} - \sum$

$G^{-1} \approx G_0^{-1} - \sum_{\text{cluster}}$

approximate $\Sigma := \sum_{\text{cluster}}$

$G^{-1} \approx G_0^{-1} - (G_{0,\text{cluster}}^{-1} - G_{\text{cluster}}^{-1})$

$G^{-1} \approx G_{\text{cluster}}^{-1} - (G_{0,\text{cluster}}^{-1} - G_0^{-1})$

$G^{-1} \approx G_{\text{cluster}}^{-1} - T$

CPT equation

Cluster Perturbation Theory in practice I

requires: **1) (numerically) exact $G_{\text{cluster}}(z, k)$**

Analytic

Exact Diagonalization (Band) Lanczos

- real frequency data
- Limited system size (~20 Fermionic orbitals)

Quantum Monte Carlo

Dynamic Density Matrix Renormalization Group

...

2) superlattice wavevector transformed T

$$T_{RR'}(\mathbf{k}) = \sum_{\mathbf{x}} e^{i\mathbf{k}\cdot\mathbf{x}} T_{Rr'} \quad \text{with} \quad r' = \mathbf{x} + R'$$

Cluster Perturbation Theory in practice II

exact limits:

$U=0$

$U/t \rightarrow \infty$

cluster size $L \rightarrow \infty$

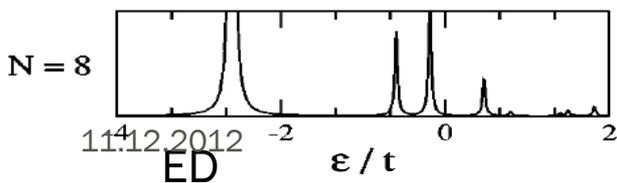
performance:

relatively simple, once $G_{\text{cluster}}(z,k)$ is obtained

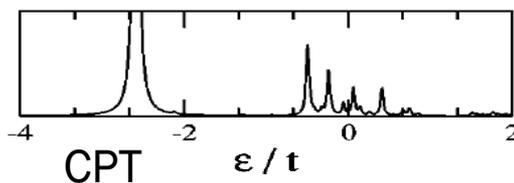
quality may depend strongly on L

improve starting point: VCA

M. Hohenadler, M. Aichhorn, W. vd Linden (2003)

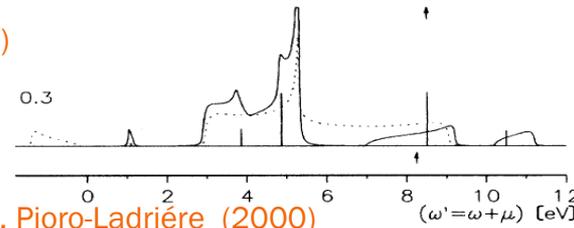


Hubbard-Holstein



C. Gros, R. Valenti (1993)

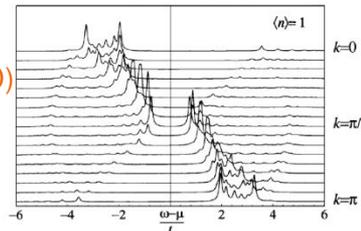
CuO₂ planes



D. Sénéchal, D. Perez, M. Pioro-Ladrière (2000)

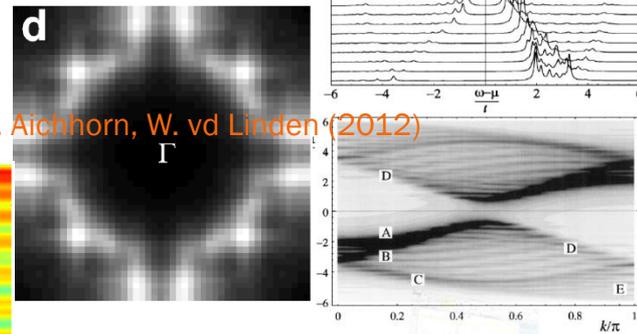
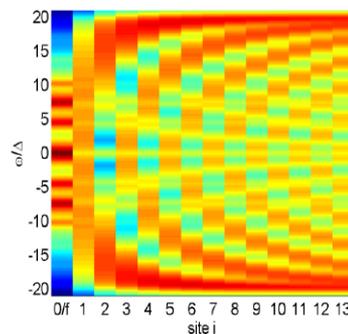
stripes in AFM

M. G. Zacher, R. Eder, E. Arrigoni, W. Hanke (2000)

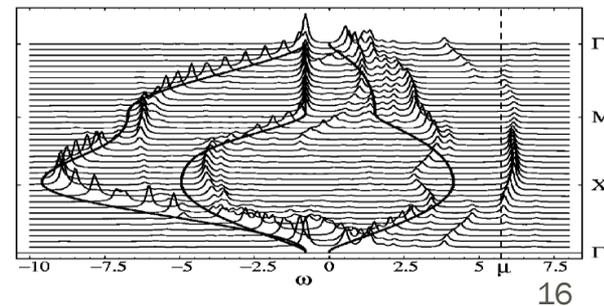


single impurity

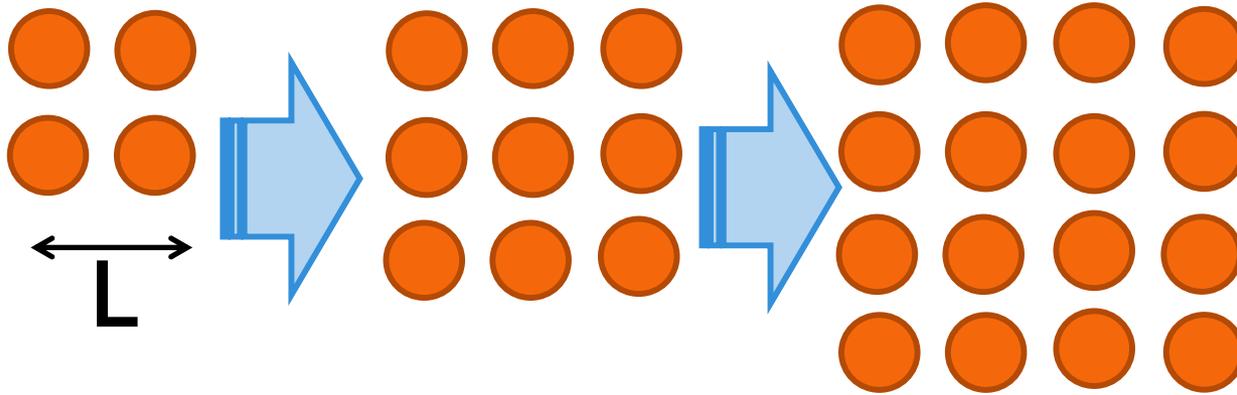
M. Nuss, E. Arrigoni, M. Aichhorn, W. vd Linden (2012)



D. Sénéchal, D. Perez, D. Plouffe (2002)



improving Cluster Perturbation Theory



CPT improves

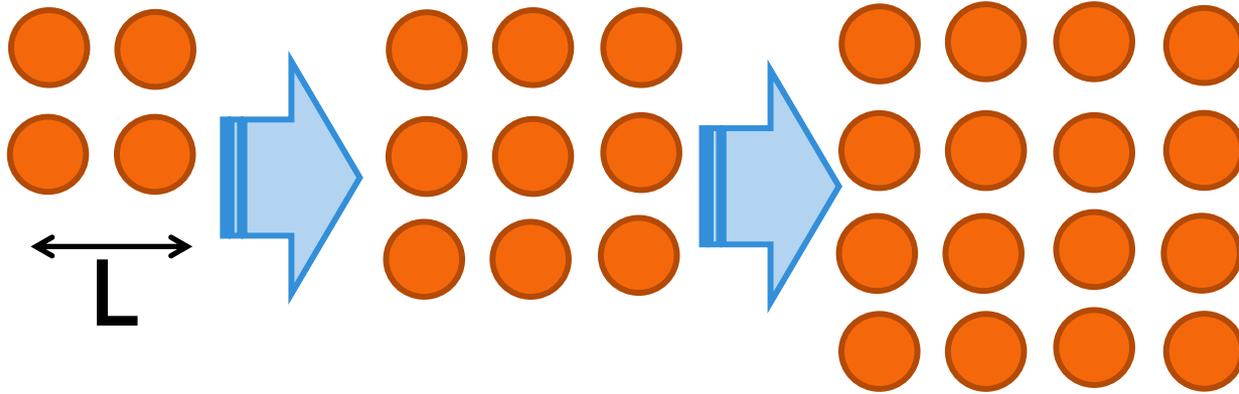
$$\Sigma_{\text{cluster}}(L)$$

already a
very fancy
tractor



but just one
steering
wheel

Variational cluster approach



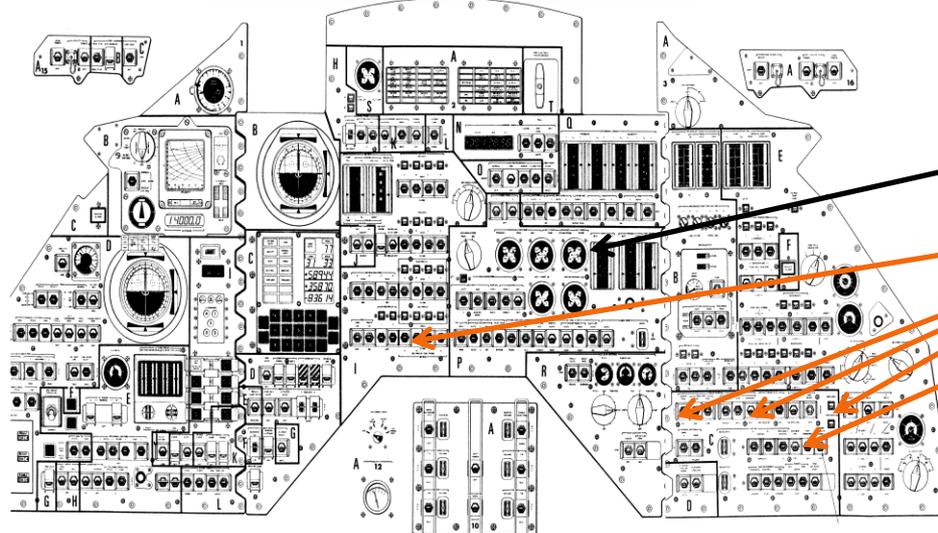
CPT improves

$$\Sigma_{\text{cluster}} (L)$$

now

if you want to land on the moon you need many more knobs to turn

APOLLO COMMAND MODULE MAIN CONTROL PANEL



$$\Sigma_{VCA} (L, \Delta_1, \Delta_2, \dots, \Delta_n)$$



Variational cluster approach

the Apollo command module however is a little bit more complicated than the fancy tractor
 VCA may be developed within the Self Energy Functional Approach M. Potthoff (2003)

$$\Omega[G] = \Phi[G] - \text{Tr} \{ (G_0^{-1} - G^{-1}) G \} + \text{Tr} \{ \ln(-G) \}$$

$$\Phi = \begin{array}{c} \text{---} \circ \\ | \\ \text{---} \circ \end{array} + \begin{array}{c} \text{---} \circ \\ \text{---} \circ \end{array} + \begin{array}{c} \text{---} \circ \\ \text{---} \circ \\ \text{---} \circ \\ \text{---} \circ \end{array} + \dots \quad \frac{\delta \Phi[G]}{\delta G} = \Sigma$$

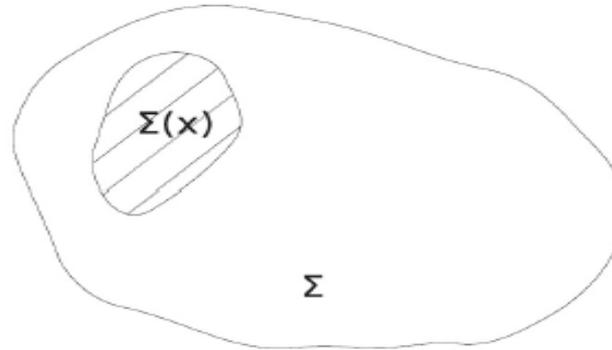
$$\frac{\delta \Omega[G]}{\delta G} = \frac{\delta \Phi[G]}{\delta G} - \frac{\text{Tr} \{ (G_0^{-1} - G^{-1}) G \}}{\delta G} + \frac{\text{Tr} \{ \ln(-G) \}}{\delta G}$$

$$= \Sigma - G_0^{-1} + G^{-1} \stackrel{!}{=} 0$$

$$\Rightarrow G^{-1} = G_0^{-1} - \Sigma$$

Variational cluster approach

limit domain of available Σ 's



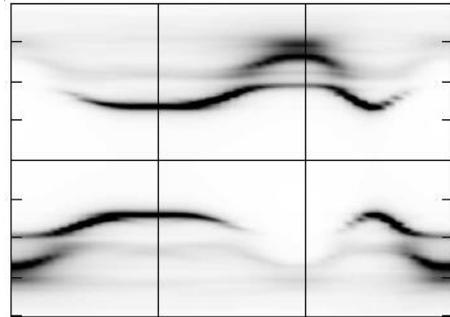
$$1) \Omega(x') = \Omega'(x') + \text{Tr} \{ \ln (-G(x')) \} - \text{Tr} \{ \ln (-G'(x')) \}$$

$$2) \nabla_{x'} \Omega(x') \stackrel{!}{=} 0$$

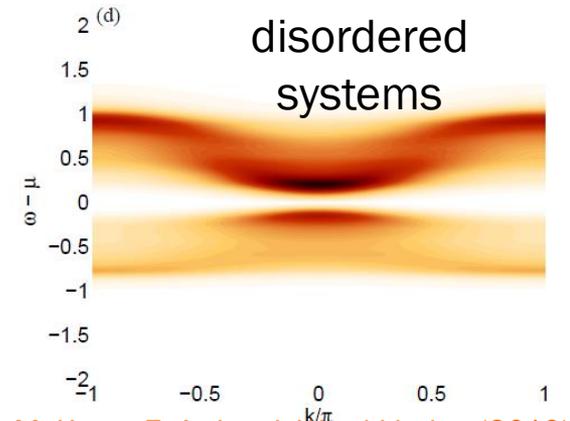
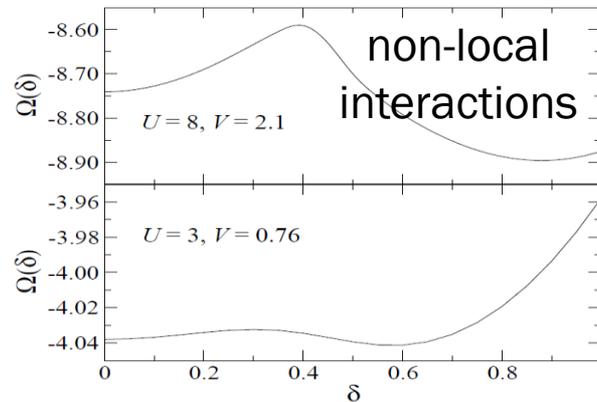
$$3) G^{-1} = G'^{-1} - T \quad \text{CPT}$$

VCA

equilibrium Variational cluster approach



spont. symmetry breaking

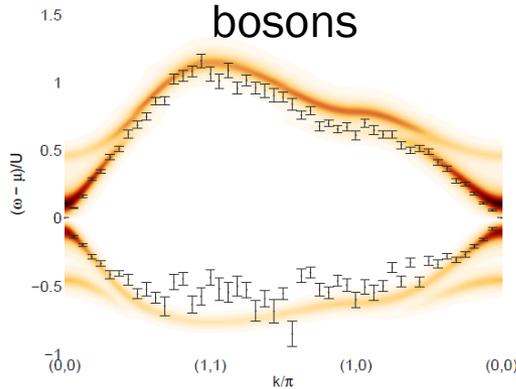


disordered systems

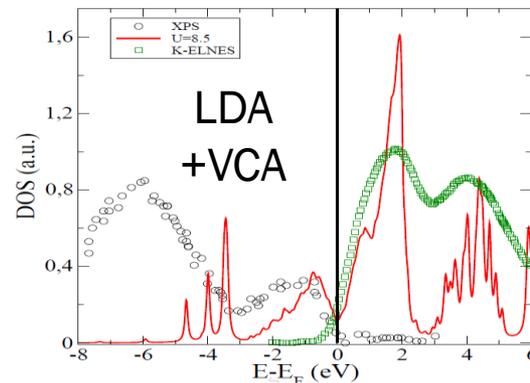
C. Dahnken, M. Aichhorn, W. Hanke, E. Arrigoni, M. Potthoff, (2004)

M. Knap, E. Arrigoni, W. vd Linden (2010)

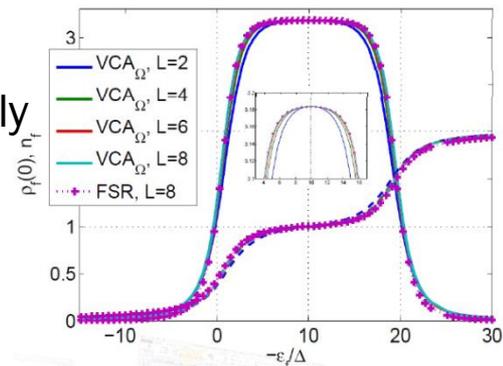
superfluid bosons



M. Aichhorn, H. G. Evertz, W. vd Linden, M. Potthoff (2004)



non-translationally invariant systems



E. Arrigoni, M. Knap, W. vd Linden (2011)

H. Allmaier, L. Chioncel, E. Arrigoni (2009)

M. Nuss, E. Arrigoni, M. Aichhorn, W. vd Linden (2012)

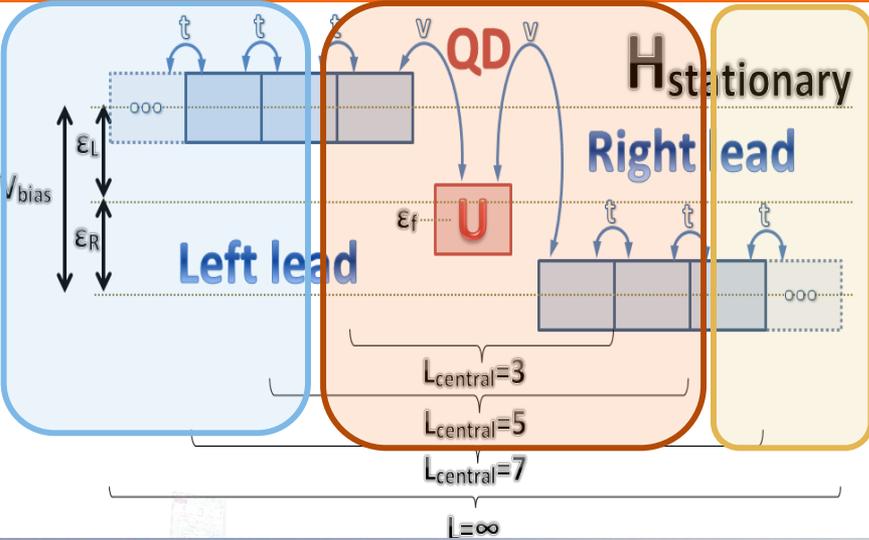
for reviews see:
11.12.2012

D.Sénéchal, arXiv.org/0806.2690 (2008)
M. Potthoff, arXiv/1108.2183 (2010)

Non-equilibrium Theory

2 - BC

Non-equilibrium Variational Cluster Approach



$$\hat{\mathcal{H}} = \hat{h} + \theta(\tau - \tau_0) \mathcal{T}$$

$\tau < \tau_0$ 3 decoupled systems: \hat{h}



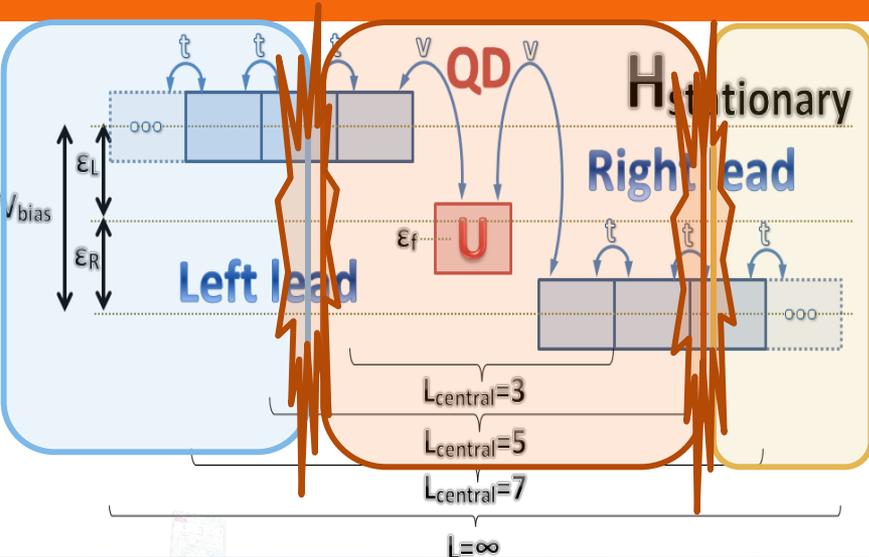
CPT time evolution

M. Balzer, M. Potthoff (2011)

VCA steady-state

M. Knap, W. von der Linden, E. Arrigoni (2011)

Non-equilibrium Variational Cluster Approach



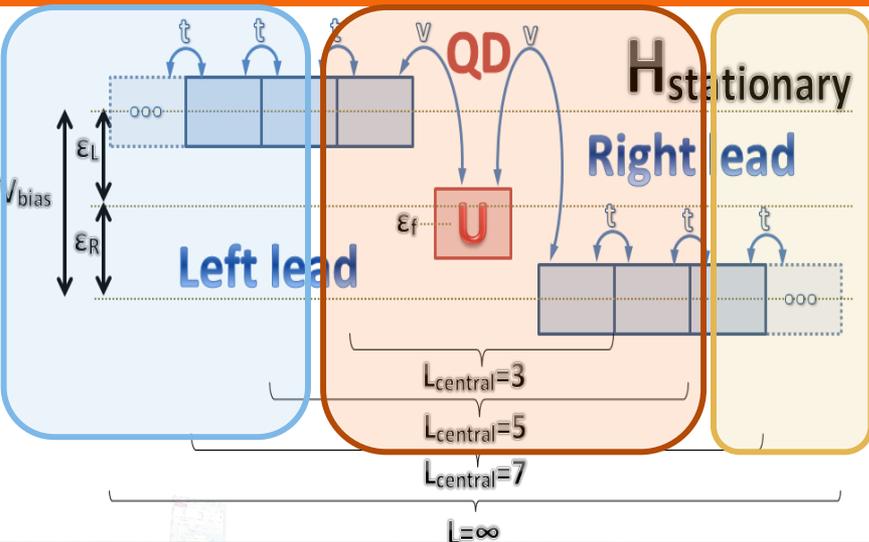
$$\hat{\mathcal{H}} = \hat{h} + \theta(\tau - \tau_0) \text{ T}$$

$\tau < \tau_0$ 3 decoupled systems: \hat{h}

@ τ_0 cpl. T switched on



Non-equilibrium Variational Cluster Approach



$$\hat{\mathcal{H}} = \hat{h} + \theta(\tau - \tau_0) \text{ [starburst symbol]}$$

$\tau < \tau_0$ 3 decoupled systems: \hat{h}

@ τ_0 cpl. [starburst symbol] T switched on

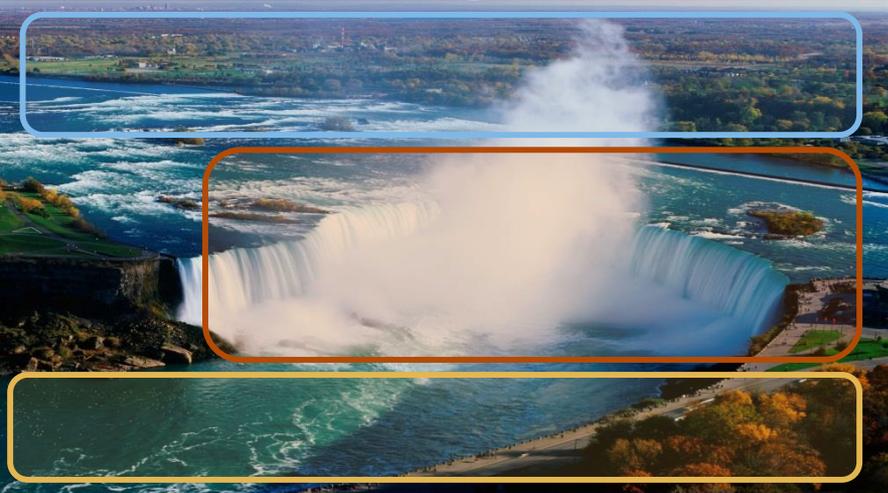
$\tau \rightarrow \infty$ steady state $\tau - \tau' \rightarrow \omega$

Keldysh

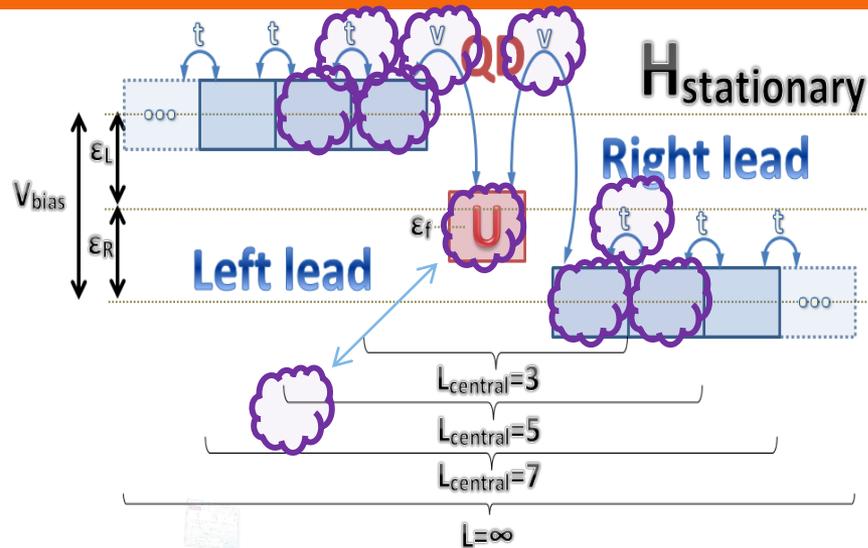
$$\tilde{G}(\omega) = \begin{pmatrix} G^R(\omega) & G^K(\omega, \mu) \\ 0 & G^A(\omega) \end{pmatrix}$$

CPT

$$\tilde{G}^{-1} = \tilde{g}^{-1} - \tilde{T}$$



Non-equilibrium Variational Cluster Approach



non unique decomposition

$$\hat{\mathcal{H}} = \hat{h} + \theta(\tau - \tau_0) \mathbb{T}$$

CPT approximation

$$\Sigma_{\hat{\mathcal{H}}} \stackrel{!}{=} \Sigma_{\hat{h}}$$

optimize initial state: VCA

$$\tau < \tau_0: \hat{h} \mapsto \hat{h} + \sum_i x_i \hat{\Delta}_i$$

$$\tau > \tau_0: \mathbb{T} \mapsto \mathbb{T} - \sum_i x_i \hat{\Delta}_i$$

flexible self-energy $\Sigma(x)$

variational principle

$$\langle \hat{\Delta}_i \rangle_{\text{initial-state}} \stackrel{!}{=} \langle \hat{\Delta}_i \rangle_{\text{steady-state}}$$

Non-equilibrium Variational Cluster Approach

$$\tau < \tau_0: \hat{h} \mapsto \hat{h} + \sum_i x_i \hat{\Delta}_i$$

$$\tau > \tau_0: T \mapsto T - \sum_i x_i \hat{\Delta}_i$$

$\langle \hat{\Delta}_i \rangle_{\text{initial-state}} \stackrel{!}{=} \langle \hat{\Delta}_i \rangle_{\text{steady-state}} = \text{self-consistent feedback}$



≈



initial reference system *as similar as possible to* **the steady-state system**

Non-equilibrium Variational Cluster Approach

full Dyson

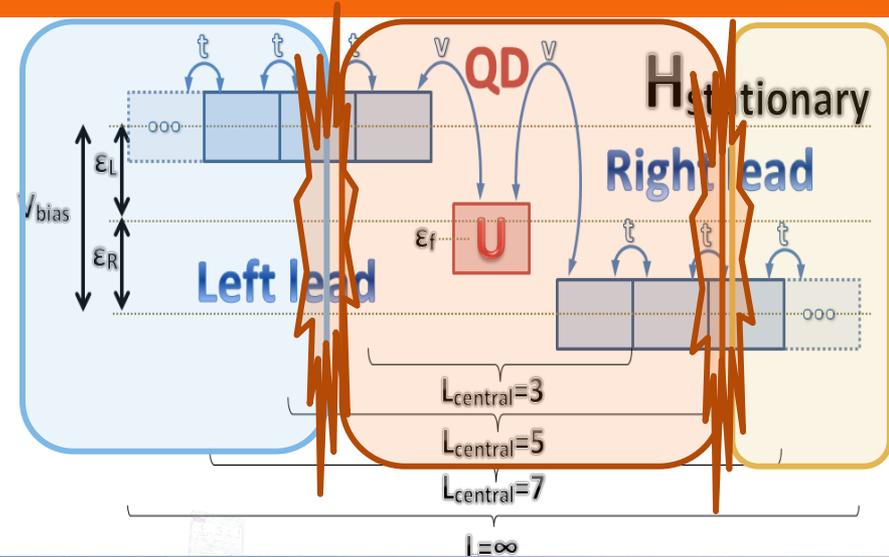
$$G = g + (T + \Delta \Sigma) G$$

CPT („Hubbard I type“)

$$G = g + (T + \Delta \Sigma) G$$

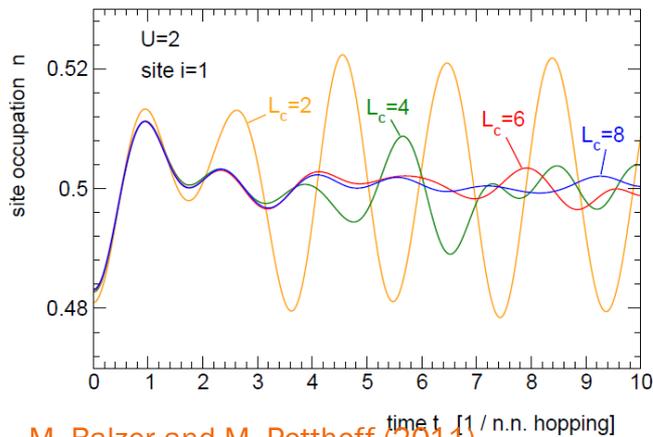
VCA

$$G = g_{\text{eff}} + (T_{\text{eff}} + \Delta \Sigma) G$$



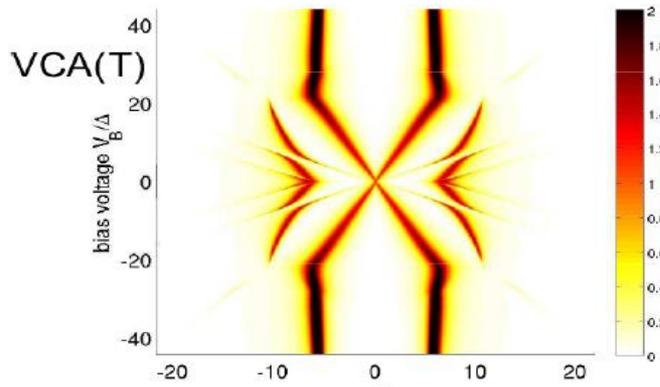
Non-equilibrium Variational Cluster Approach

Time evolution by CPT



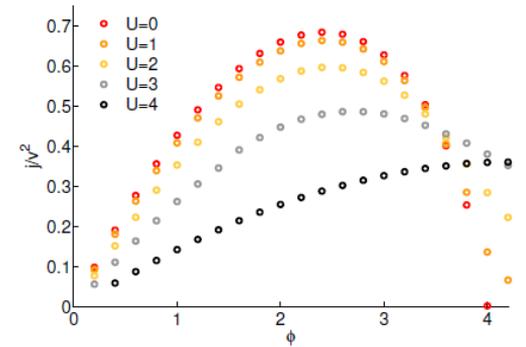
M. Balzer and M. Potthoff (2011)

Steady state of quantum dots



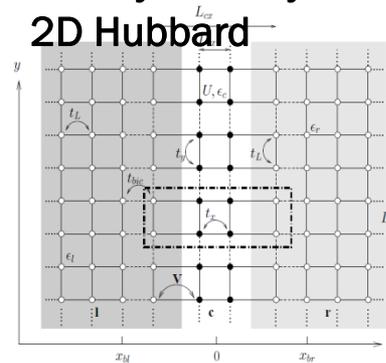
M. Nuss, C. Heil, M. Ganahl, M. Knap, H. G Evertz, E. Arrigoni, W. vd Linden (2012)

non equilibrium DMFT



E. Arrigoni, M. Knap, W. vd Linden (2012)

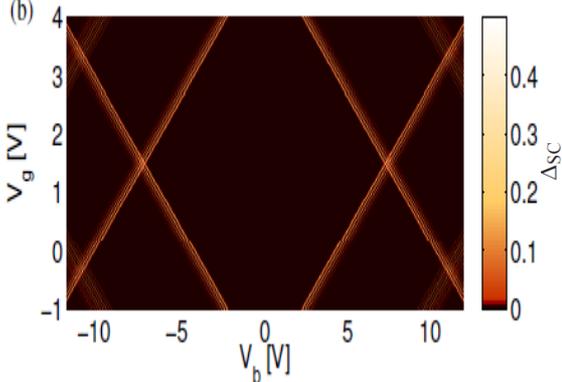
Steady state by VCA, 2D Hubbard



M. Knap, W. vd Linden, E. Arrigoni (2011)

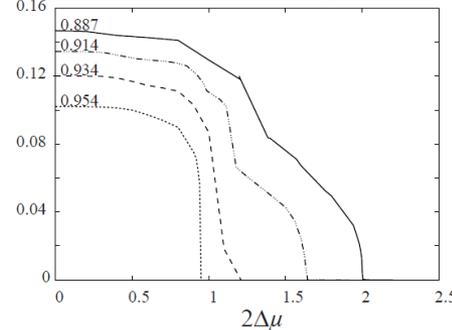
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Phonons in molecular devices



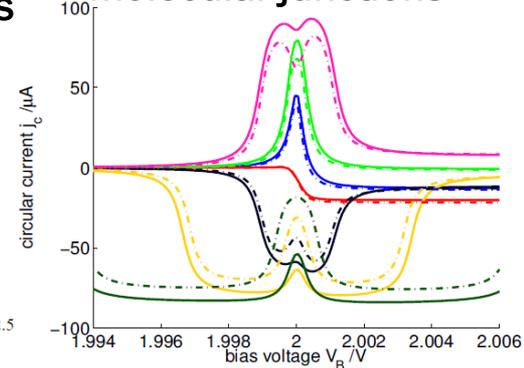
M. Knap, E. Arrigoni, W. von der Linden (2012)

Superconducting layers



A. Fulterer, E. Arrigoni (2012)

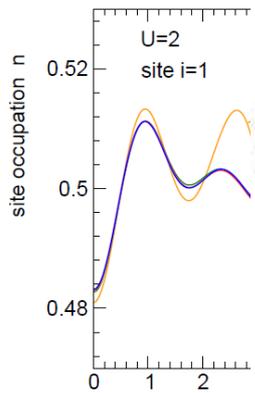
Magnetic effects in molecular junctions



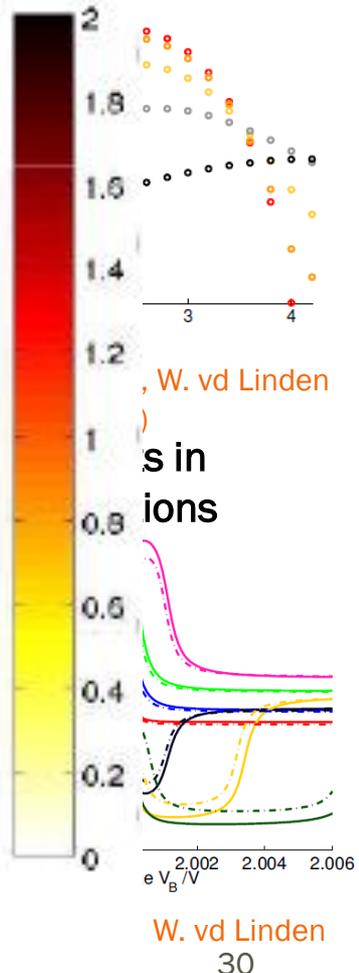
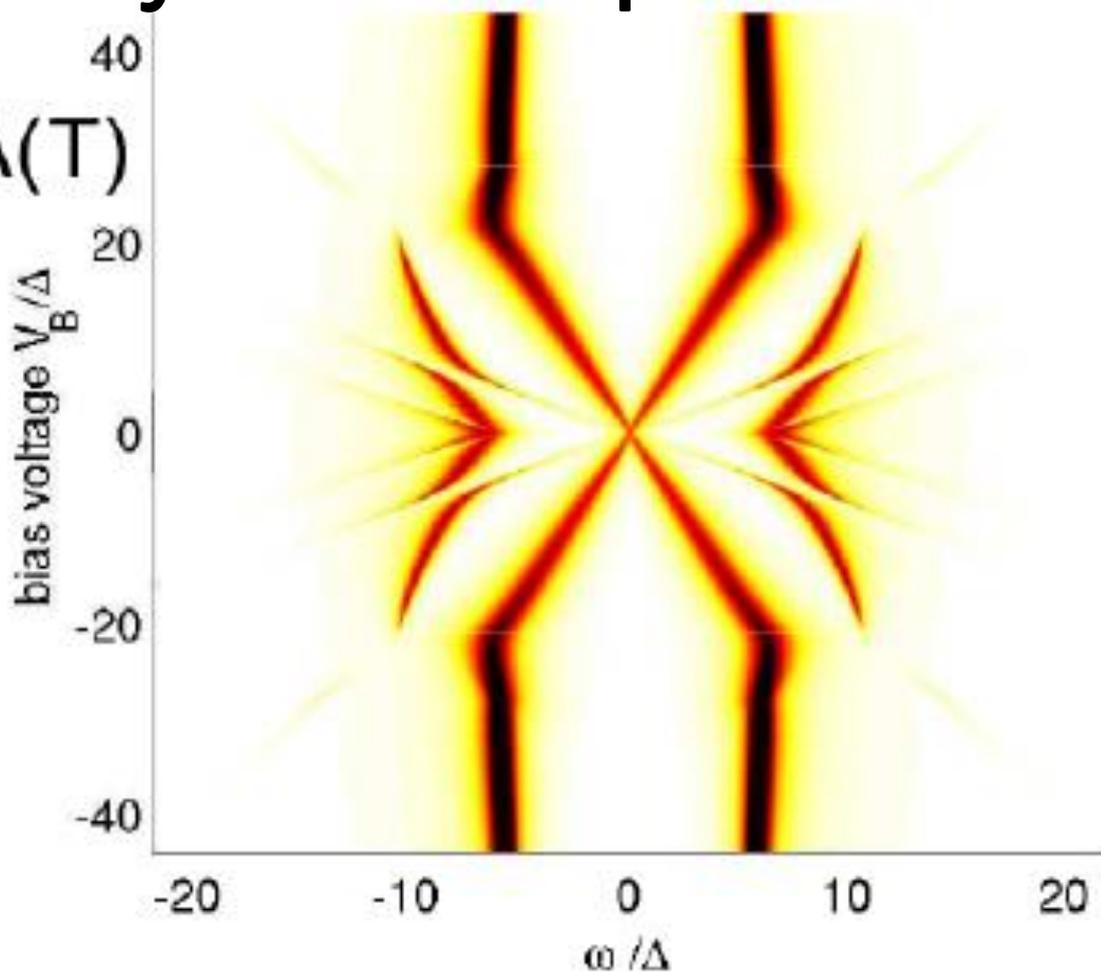
M. Nuss, E. Arrigoni, W. vd Linden (2012)

Non-equilibrium Variational Cluster Approach

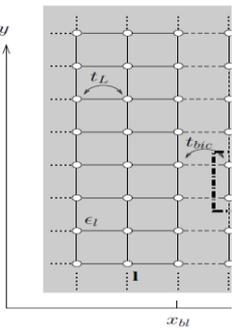
Steady state of quantum dots



VCA(T)



M. Balzer and M
Steady state



M. Knap, W. vd L

11.12.2012

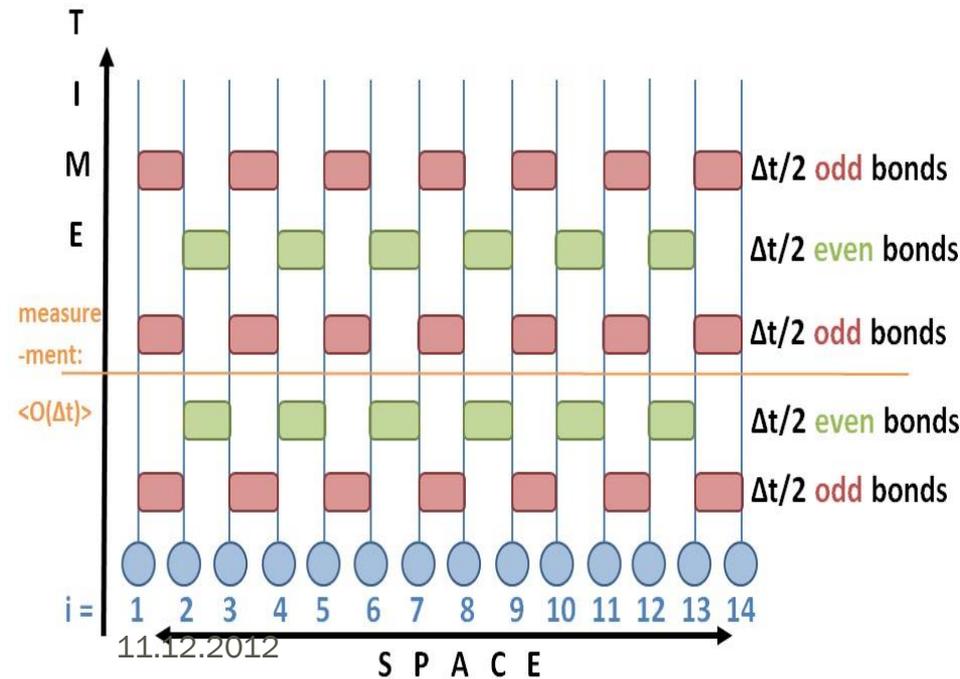
M. Nuss, C. Heil, M. Ganahl, M. Knap, H. G Evertz, E. Arrigoni, W. vd Linden

Matrix Product State time evolution



Time evolution with Matrix Product States

$$\begin{aligned}
 |\Psi\rangle &= \sum_{\{s_1, s_2, \dots, s_L\}} c_{s_1, s_2, \dots, s_L} |s_1, s_2, \dots, s_L\rangle \\
 &= \sum_{\{s_1, \dots, s_L\}} \sum_{\{\alpha_1, \dots, \alpha_L\}} A_{\alpha_1}^{[1]s_1} A_{\alpha_1 \alpha_2}^{[2]s_2} \dots A_{\alpha_{L-2} \alpha_{L-1}}^{[L-1]s_{L-1}} A_{\alpha_{L-1}}^{[L]s_L} |s_1, \dots, s_L\rangle
 \end{aligned}$$



S. R. White (1993)

$$1) \text{ DMRG}(H_0(t_0)) \longrightarrow \Psi_0$$

$$2) \text{ @ } t_x \text{ quench: } H_0 \longrightarrow H_>$$

$$3) \text{ TEBD}(\Psi_0) \longrightarrow \Psi(t)$$

G. Vidal (2004)

Time evolution with Matrix Product States

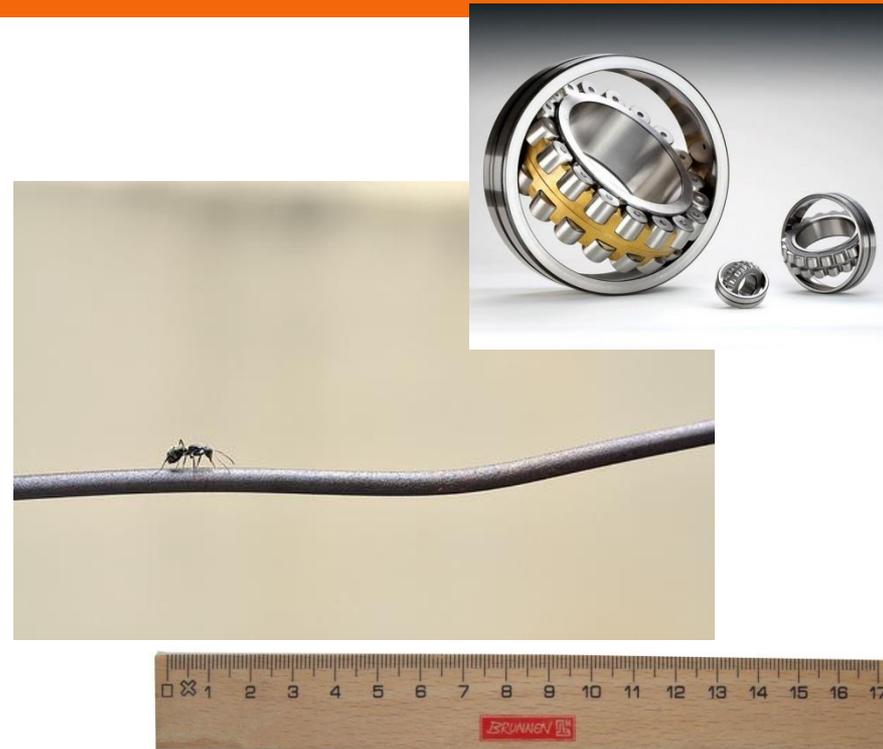
quasi exact / high
precision

most powerful in **1D**

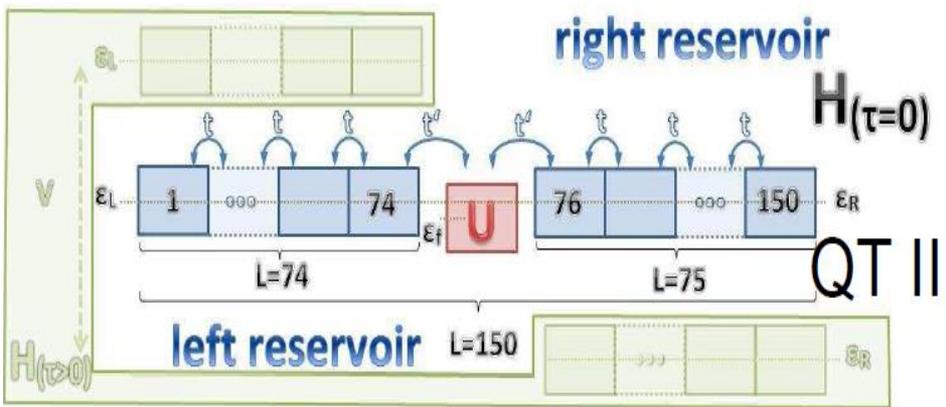
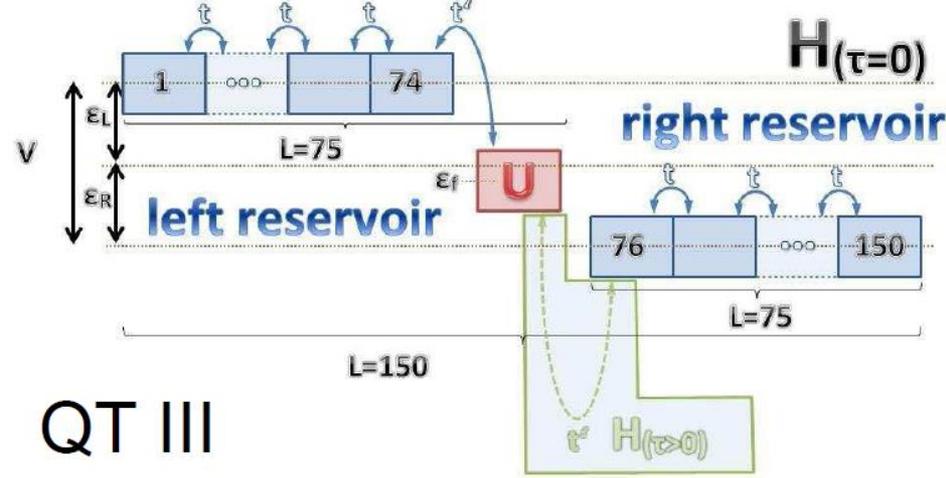
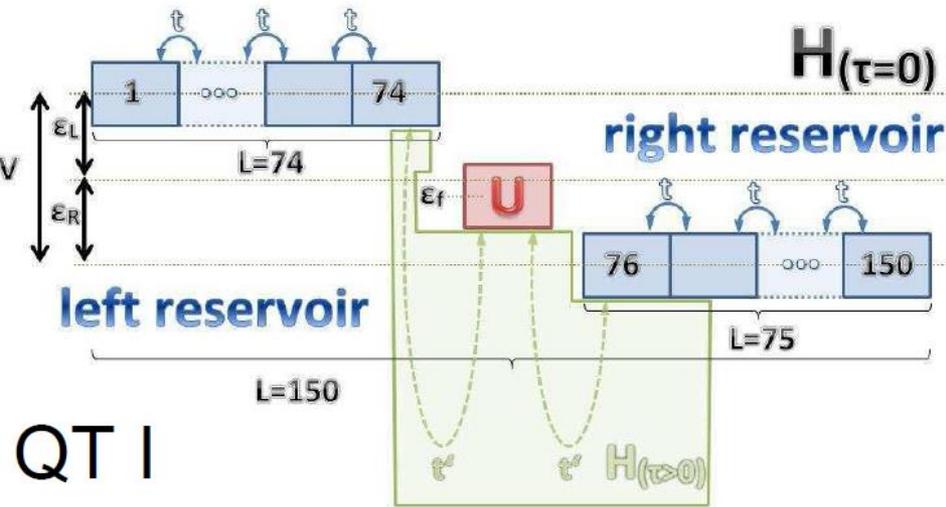
limited system **size**

limited simulation **time**

11.12.2012



Time evolution after a quantum quench



comparison of short time quench dynamics + quasi-steady-state

M. Nuss, M. Ganahl, H. G. Evertz, E. Arrighi, W. von der Linden (in preparation)

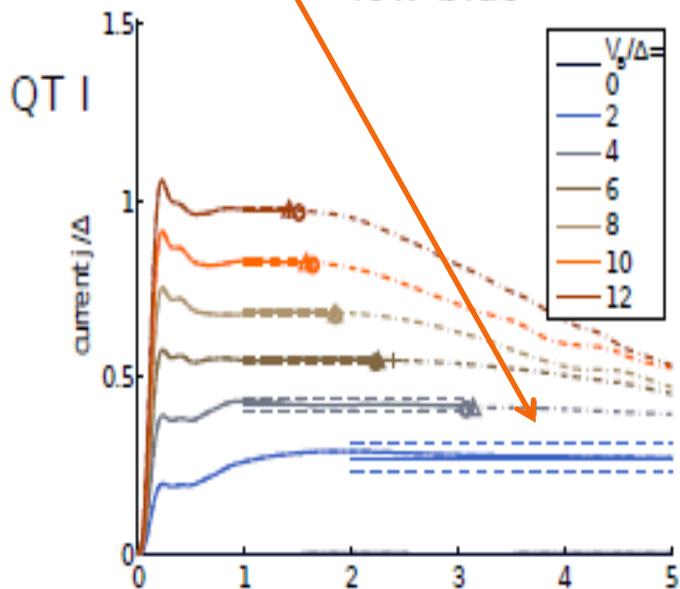
Can the steady state be reached?

entanglement

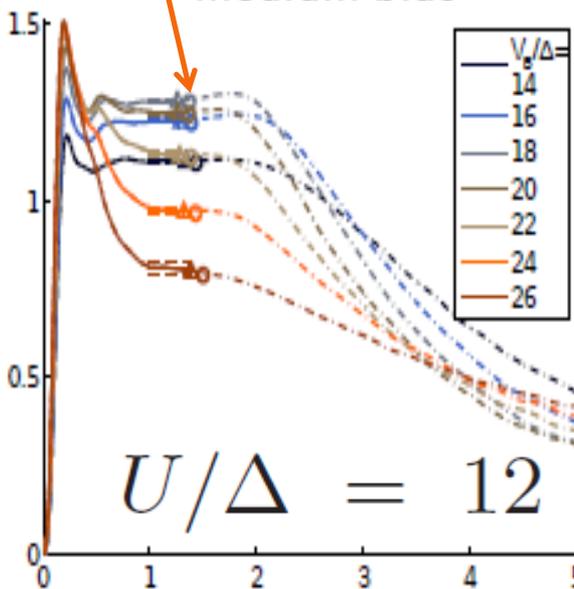


steady-state ✓

low bias

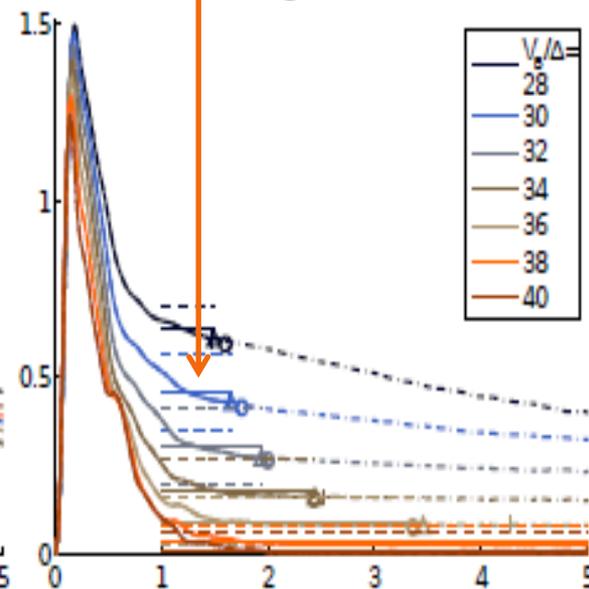


medium bias



steady-state ✗

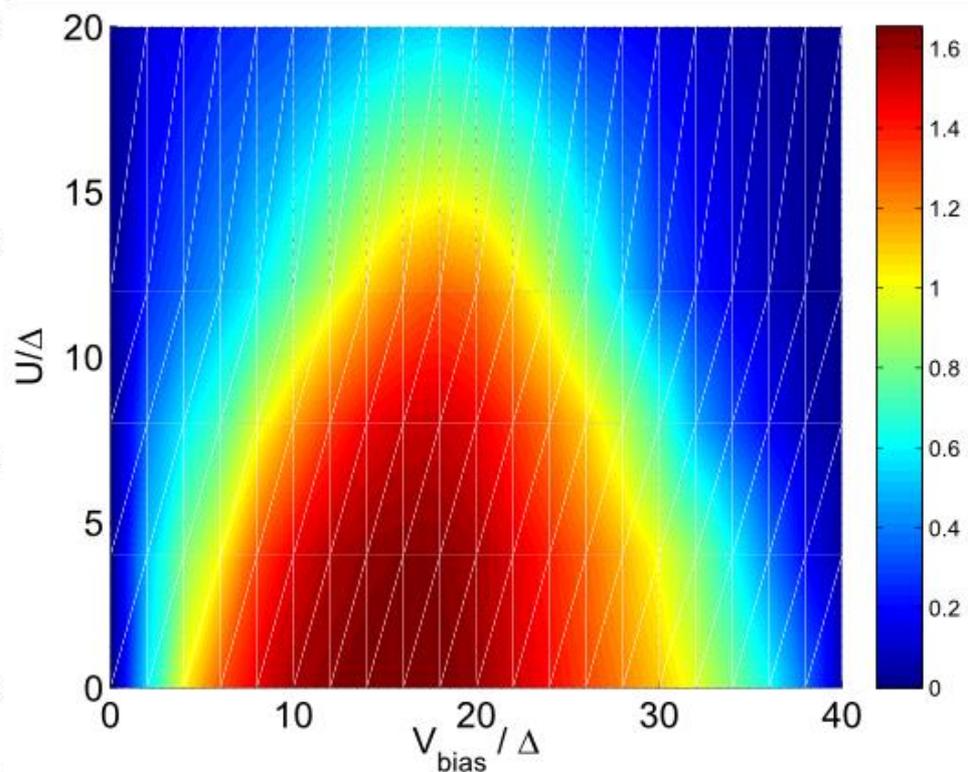
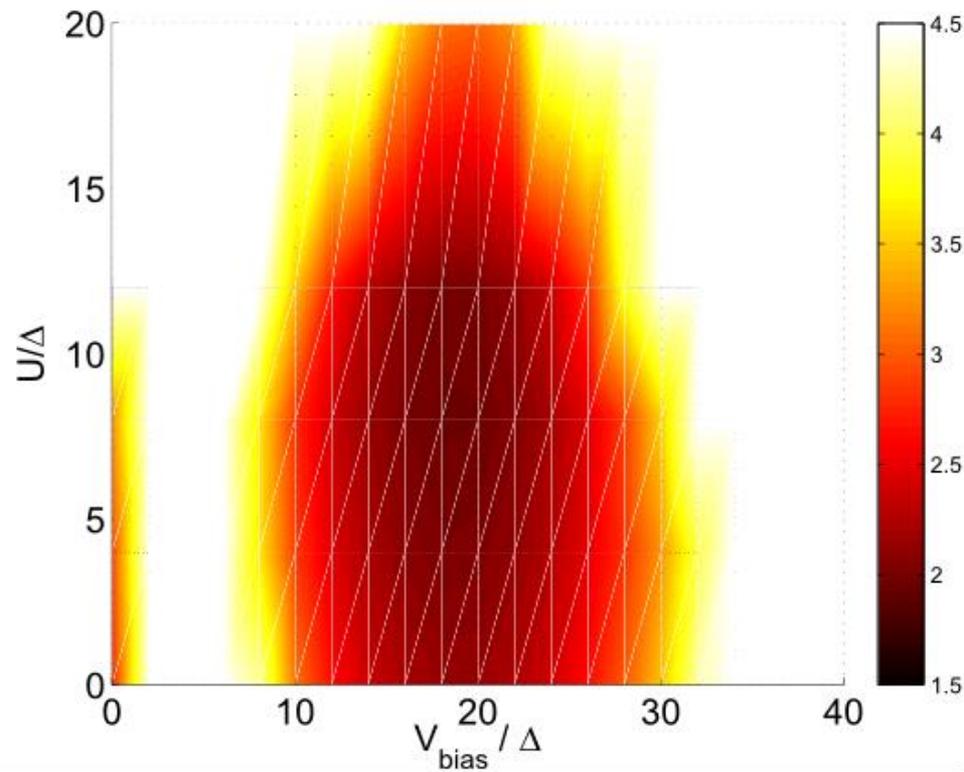
high bias



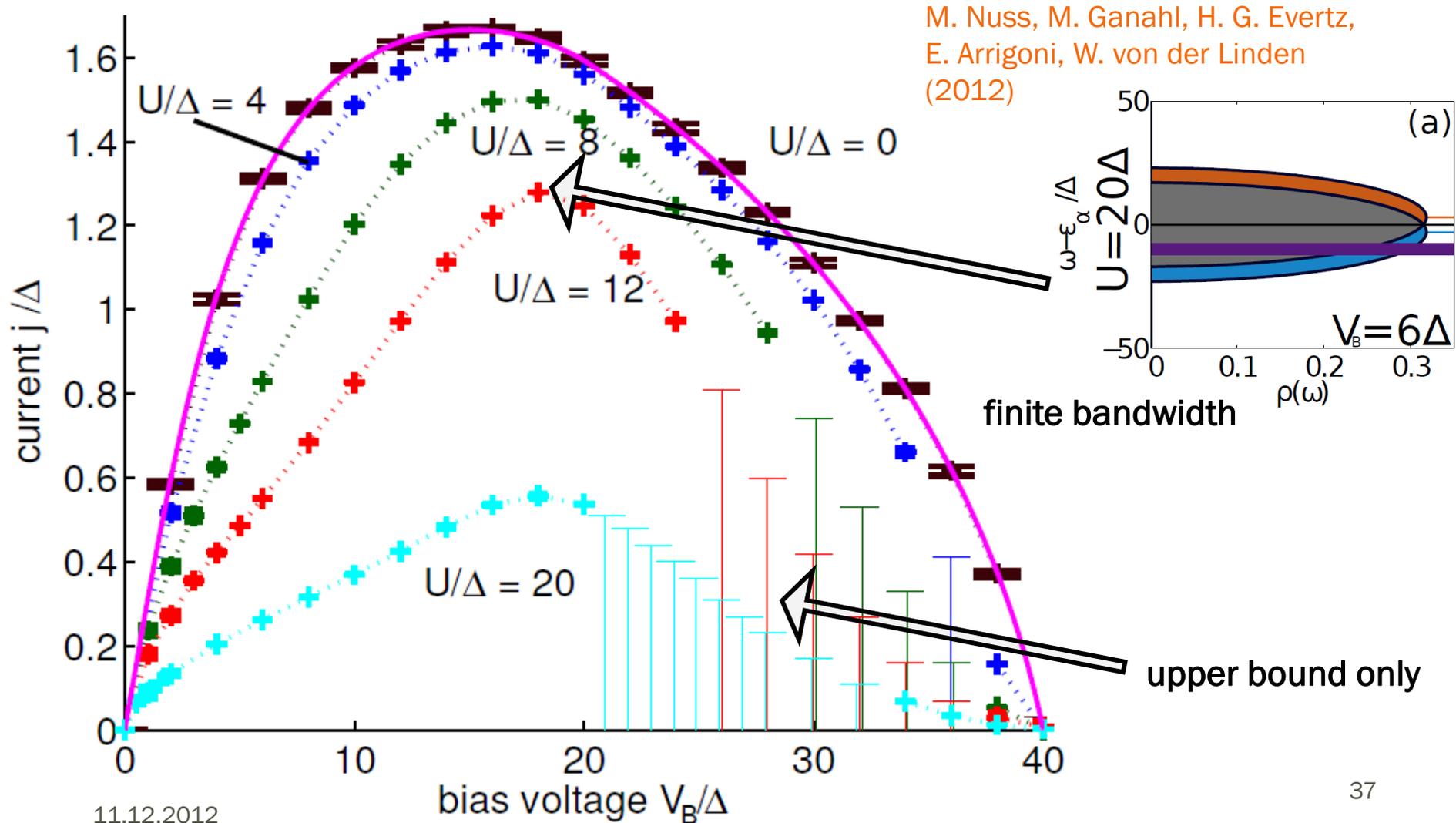
entanglement



current



DMRG+TEBD current-voltage characteristics



Steady-state properties of a quantum dot

∞ 4 ∞

equilibrium cluster theory for SIAM

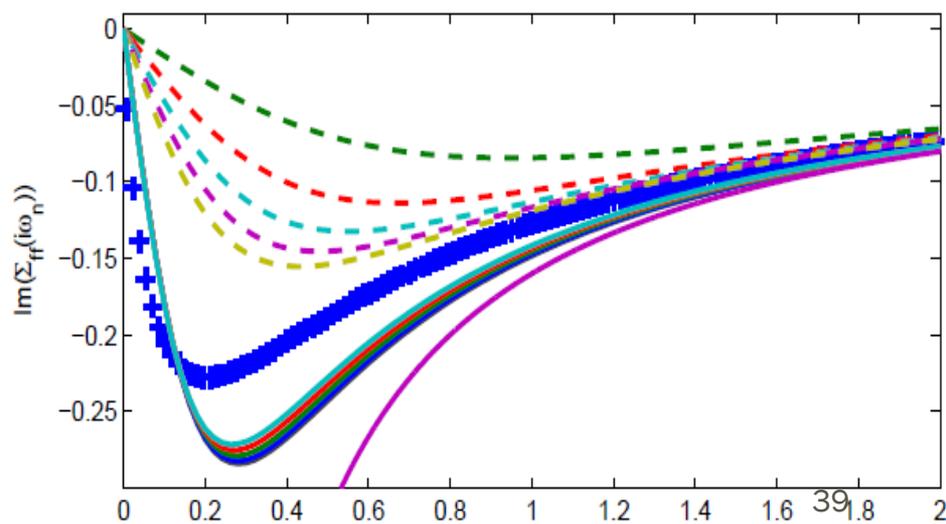
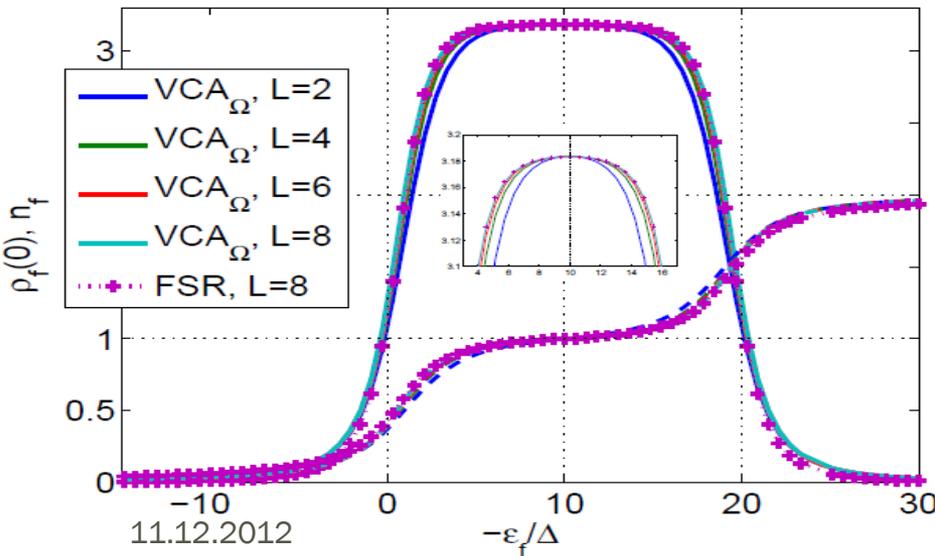
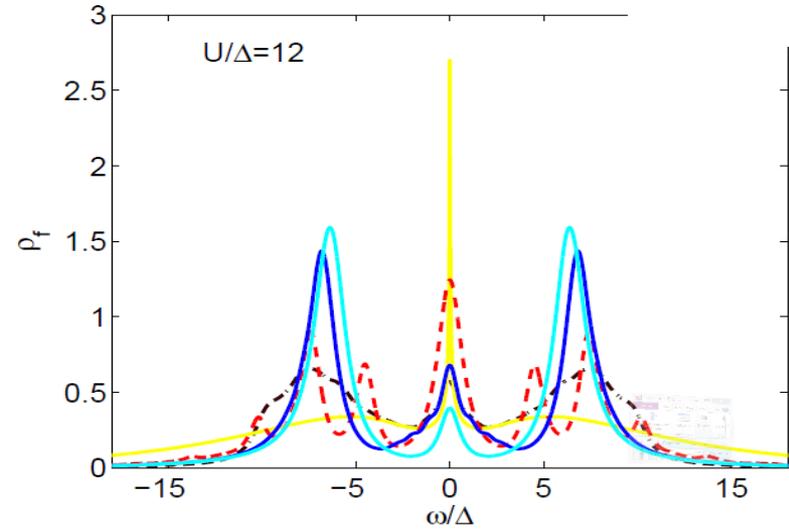
much more accurate methods, c.f. NRG, ...

exponentially large Kondo singlet (Bethe Ansatz)

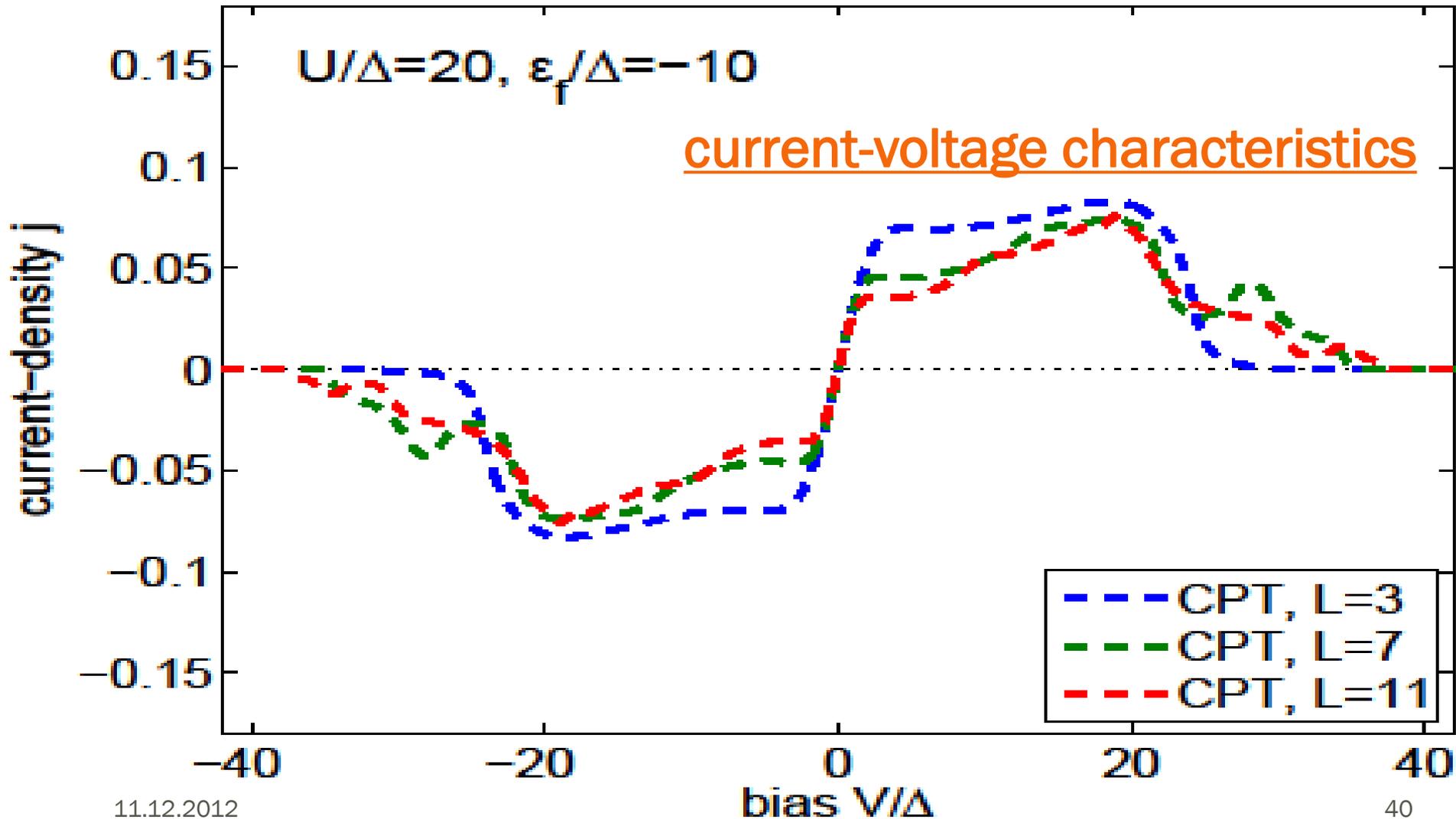
$$T_K = \sqrt{\frac{\Delta U}{2}} e^{-\gamma \frac{\pi}{8\Delta} U}, \quad \gamma = 1 \quad \text{(VCA)=0.651}$$

Friedel sum rule

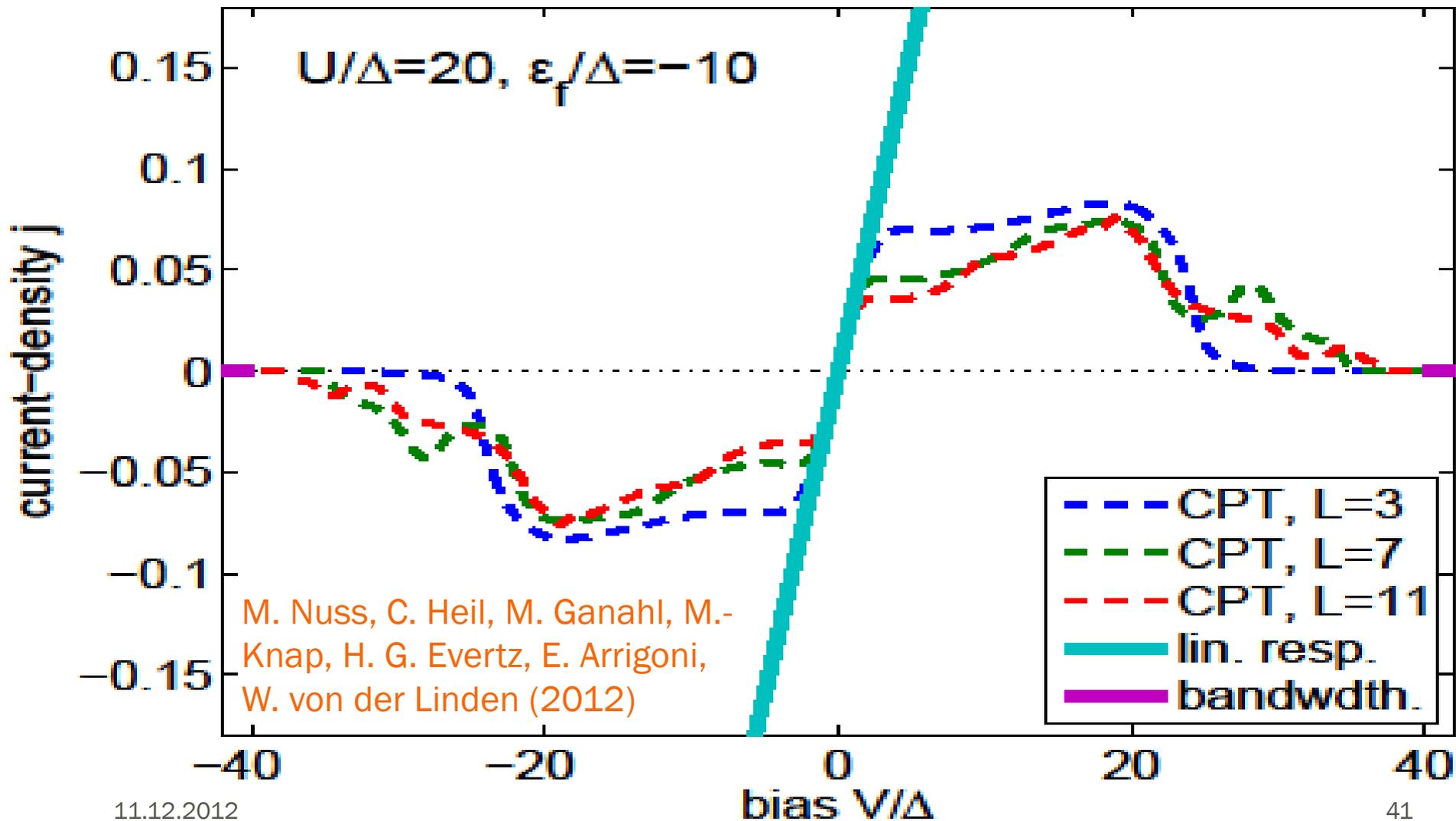
$$\rho_f(0) = \frac{1}{\pi\Delta} \sin^2\left(\frac{\pi\langle n^f \rangle}{2}\right)$$



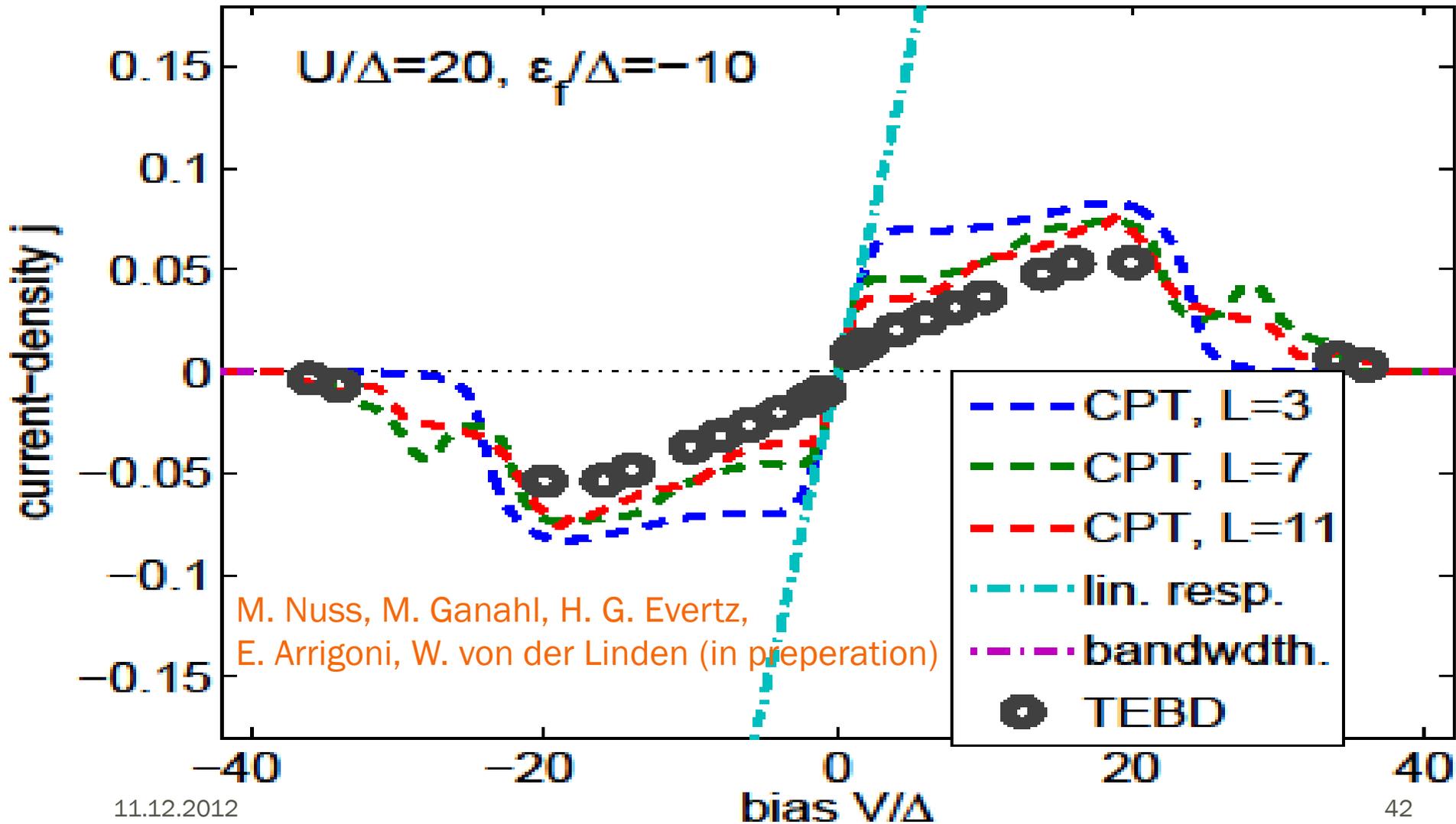
non-equilibrium CPT



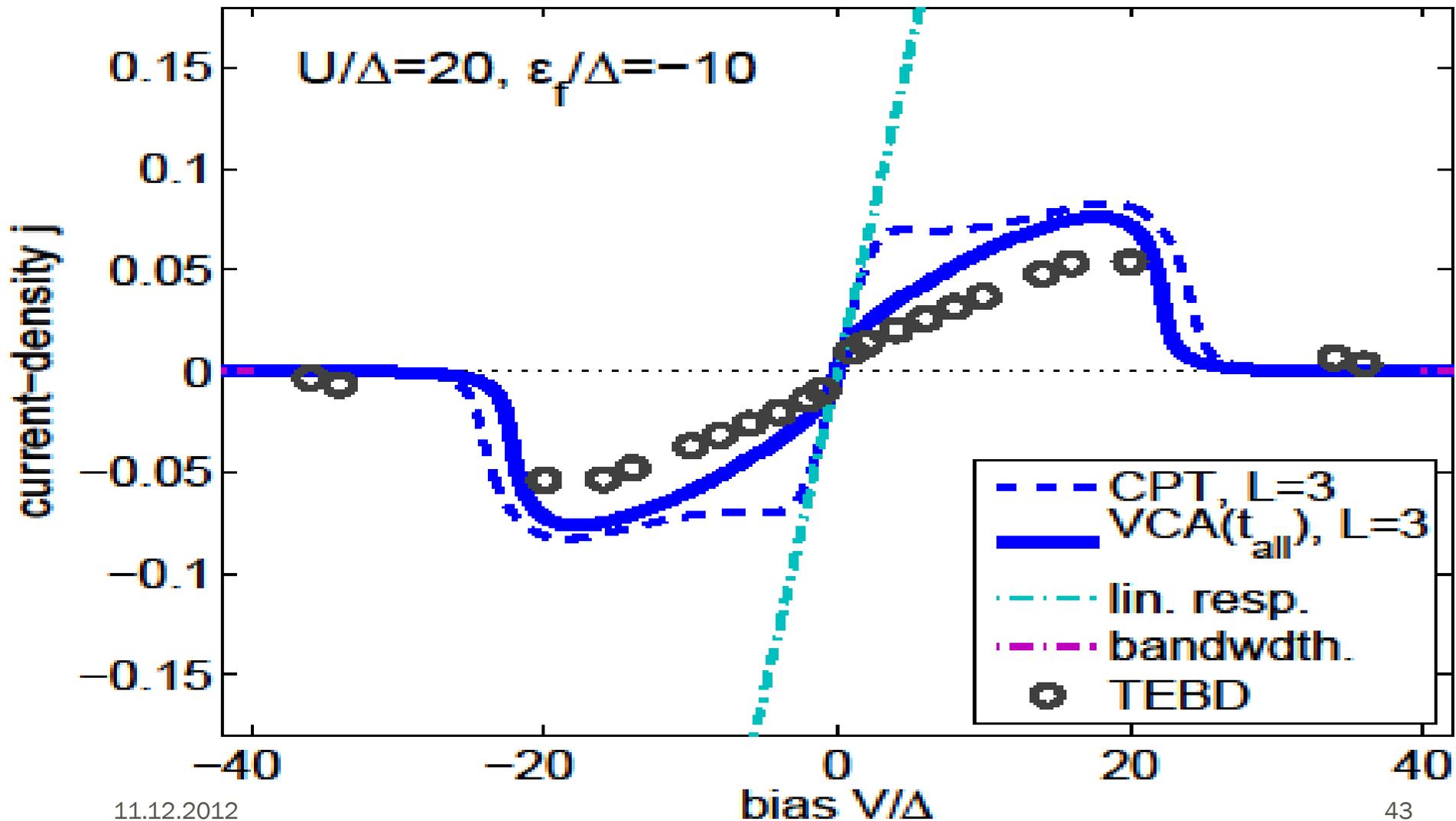
exact limits



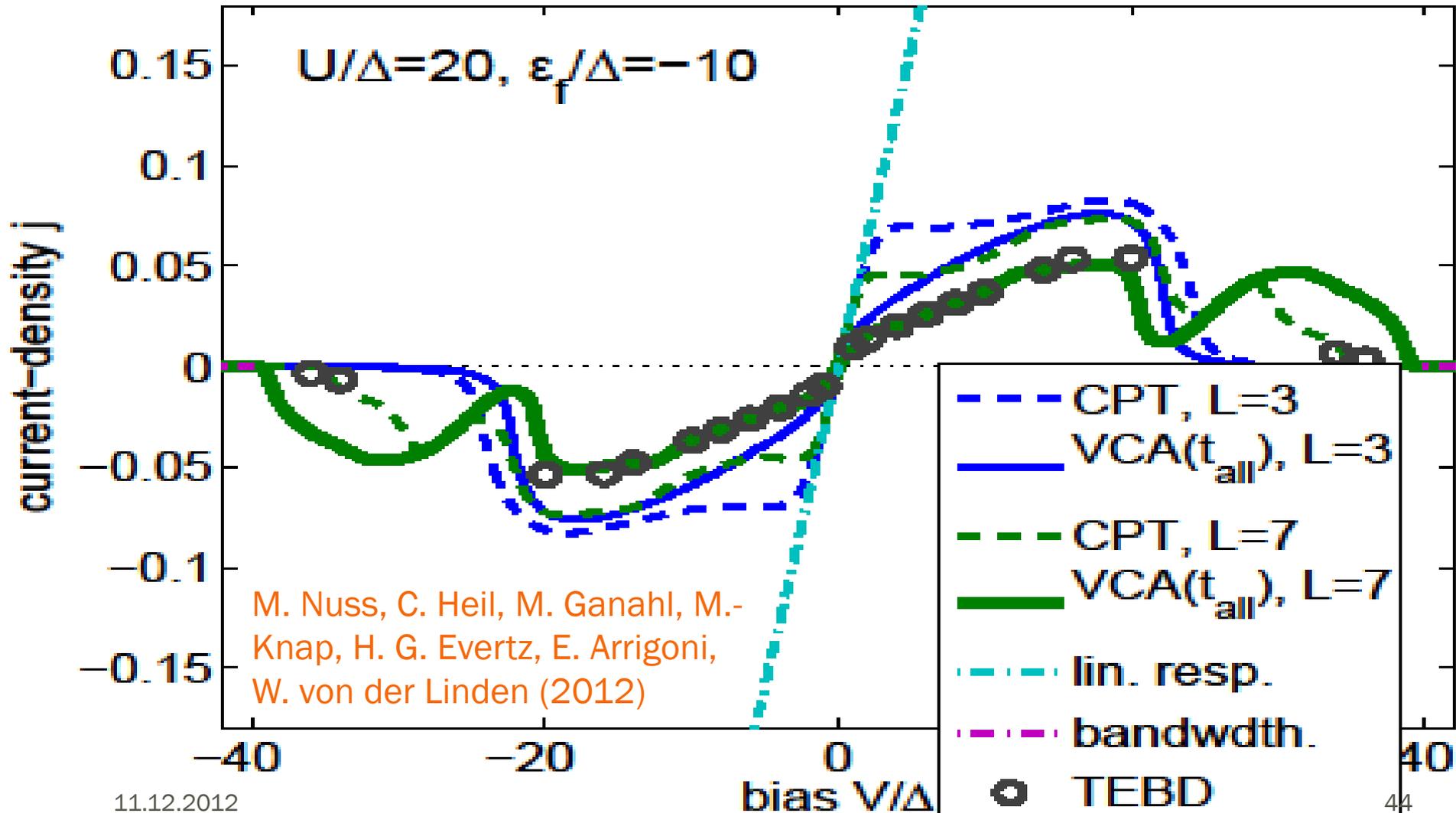
quasi-exact DMRG+TEBD



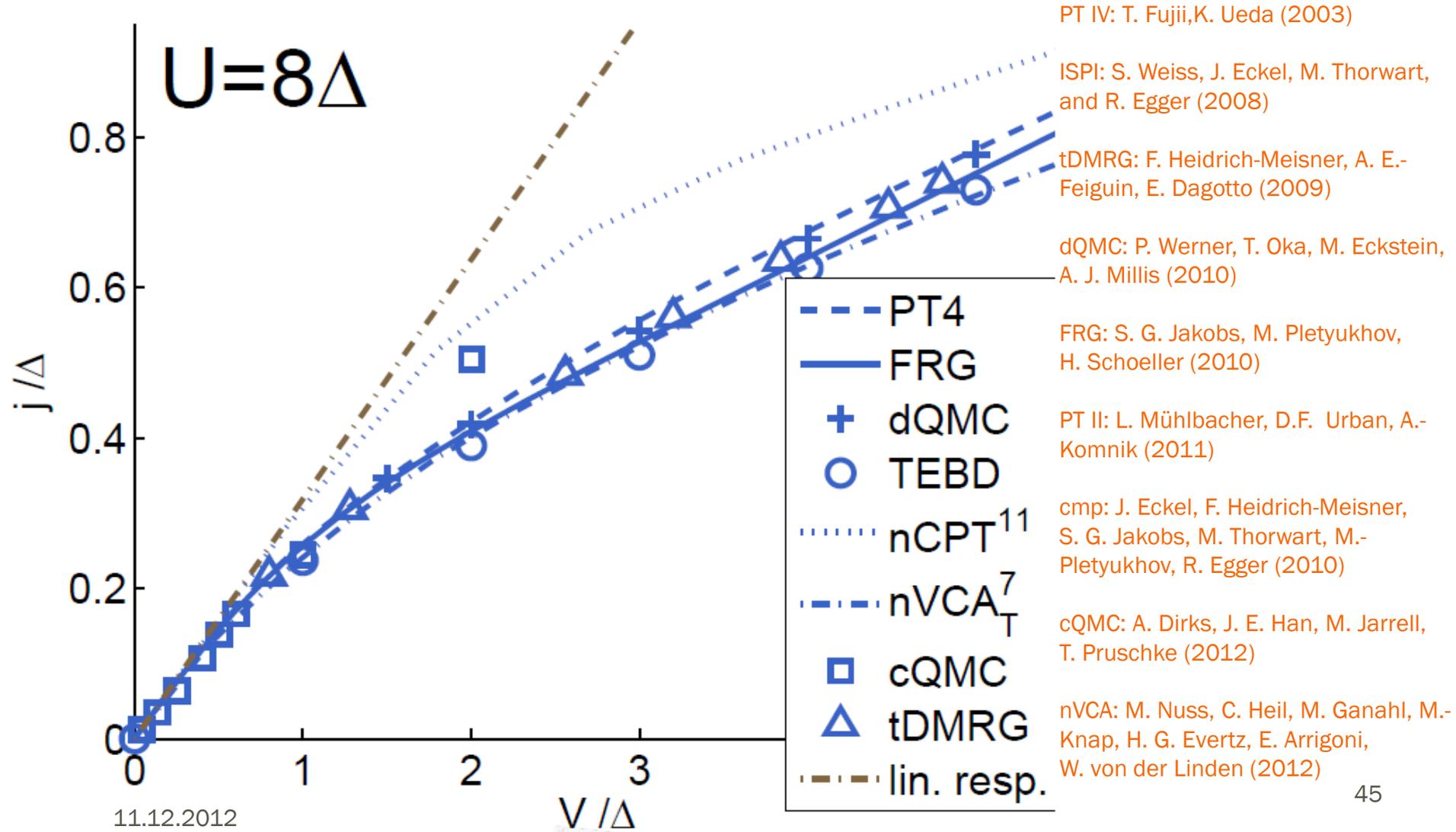
non-equilibrium VCA



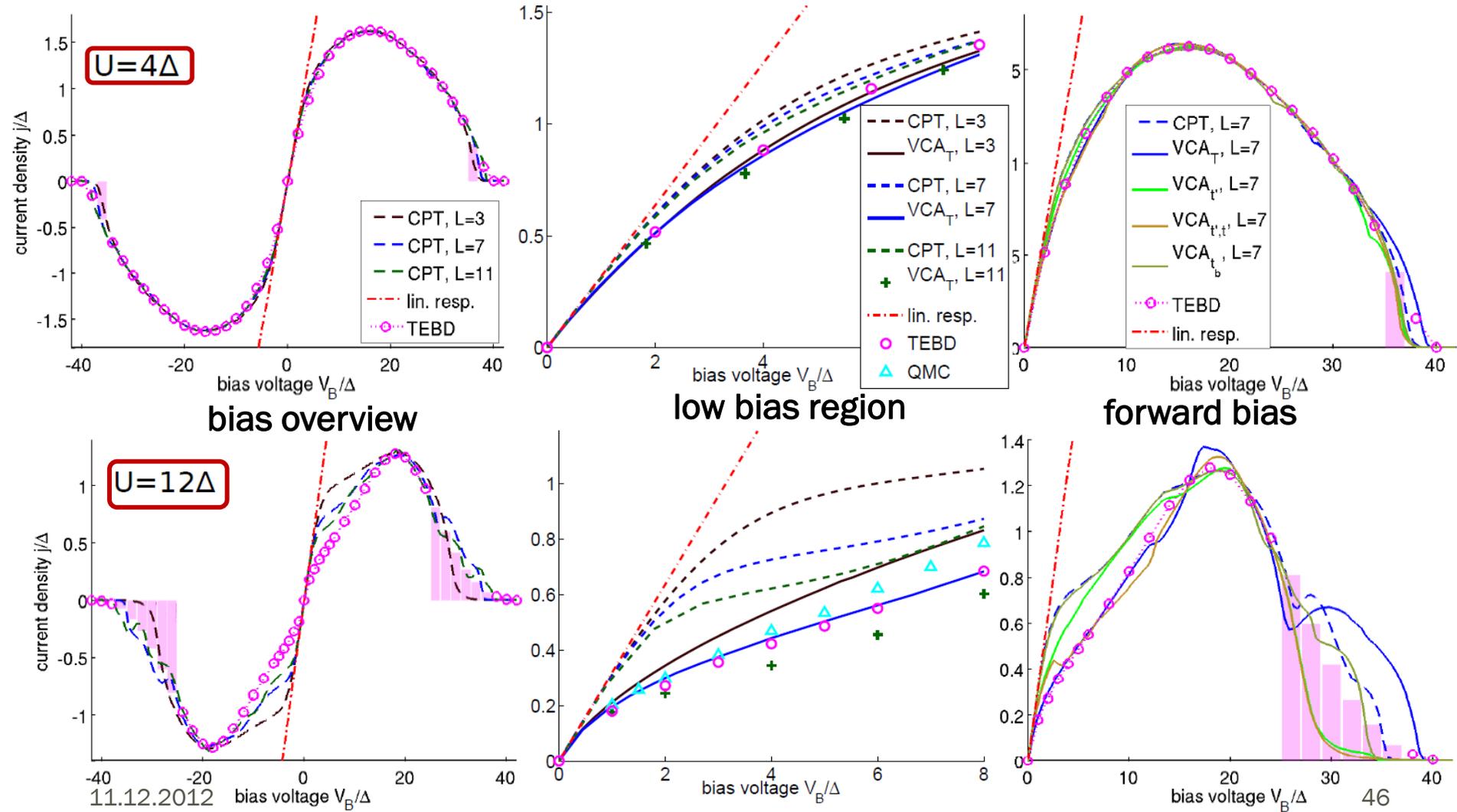
non-equilibrium VCA



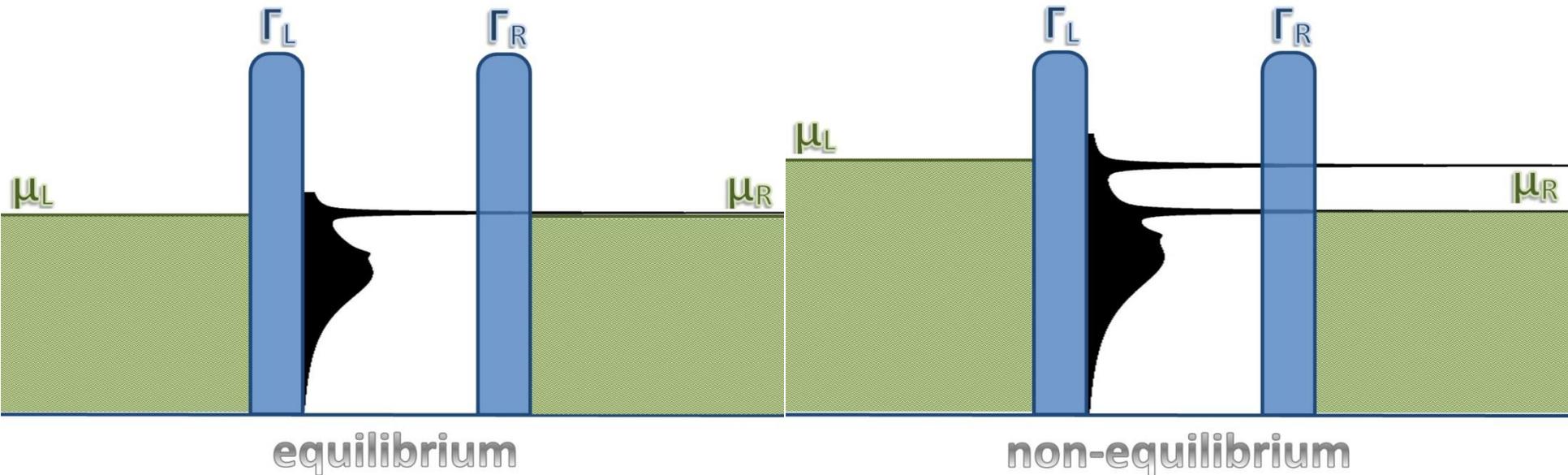
comparison in low bias regime



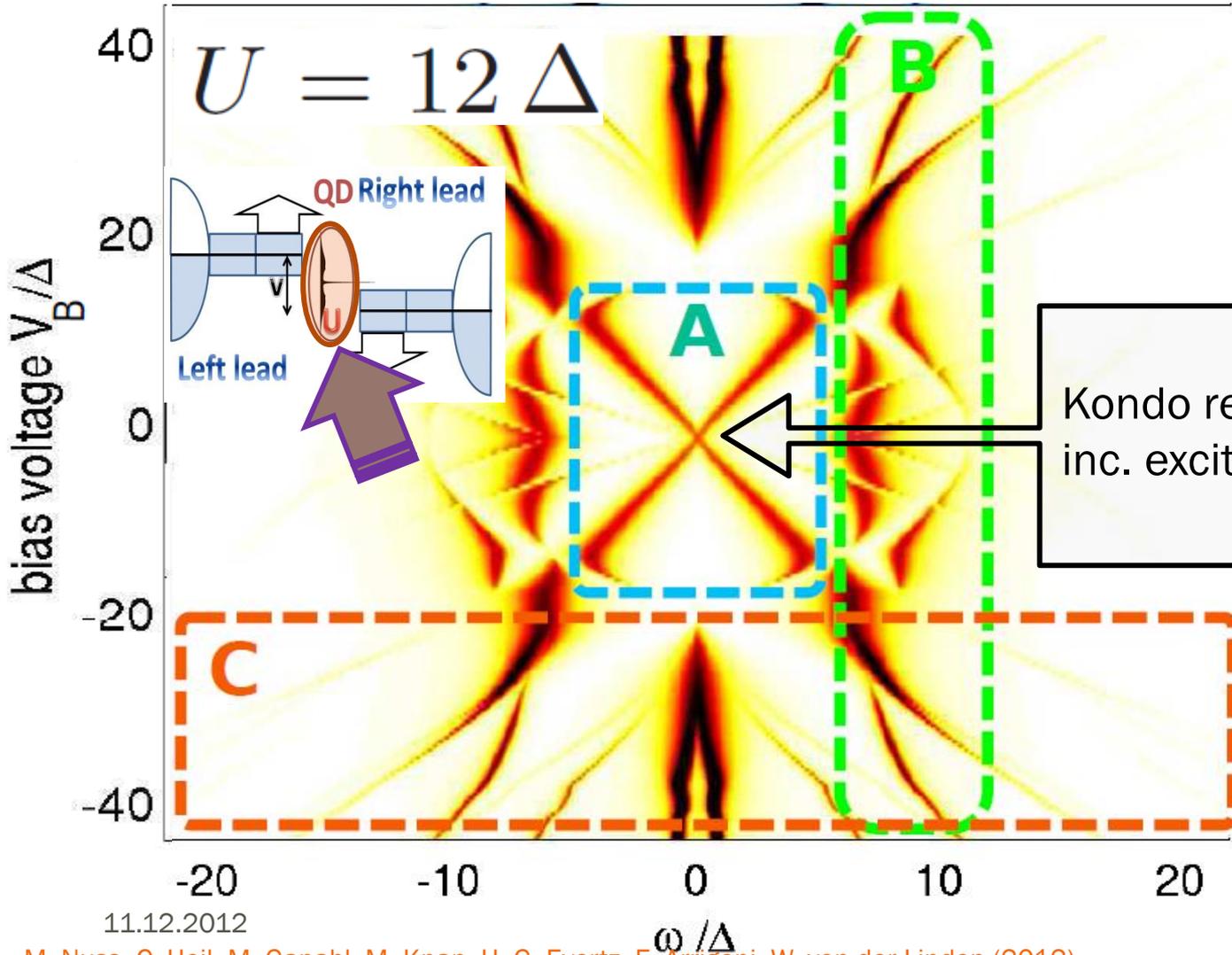
current voltage performance



non-equilibrium local density of states



non-equilibrium local density of states



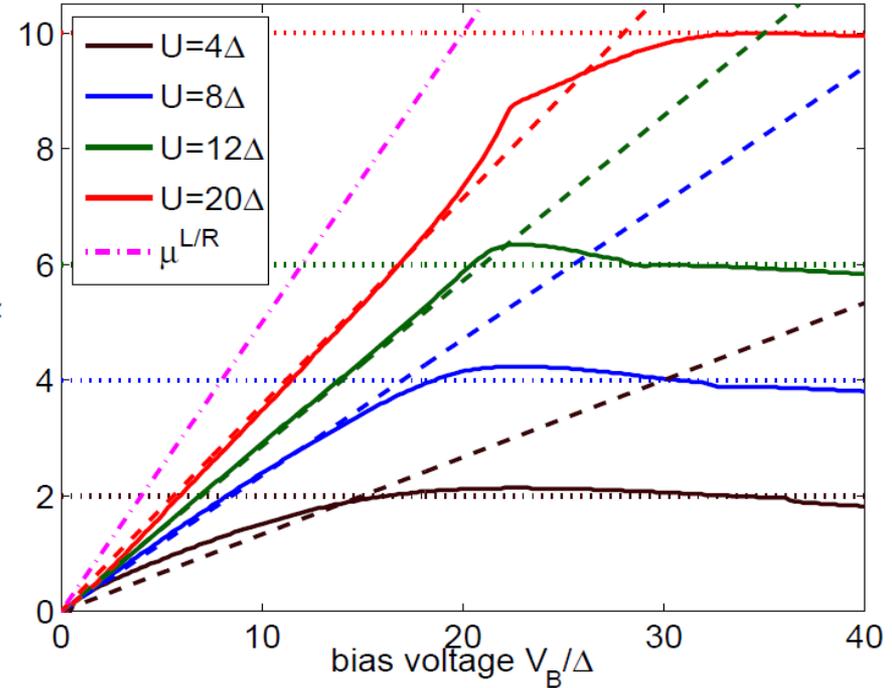
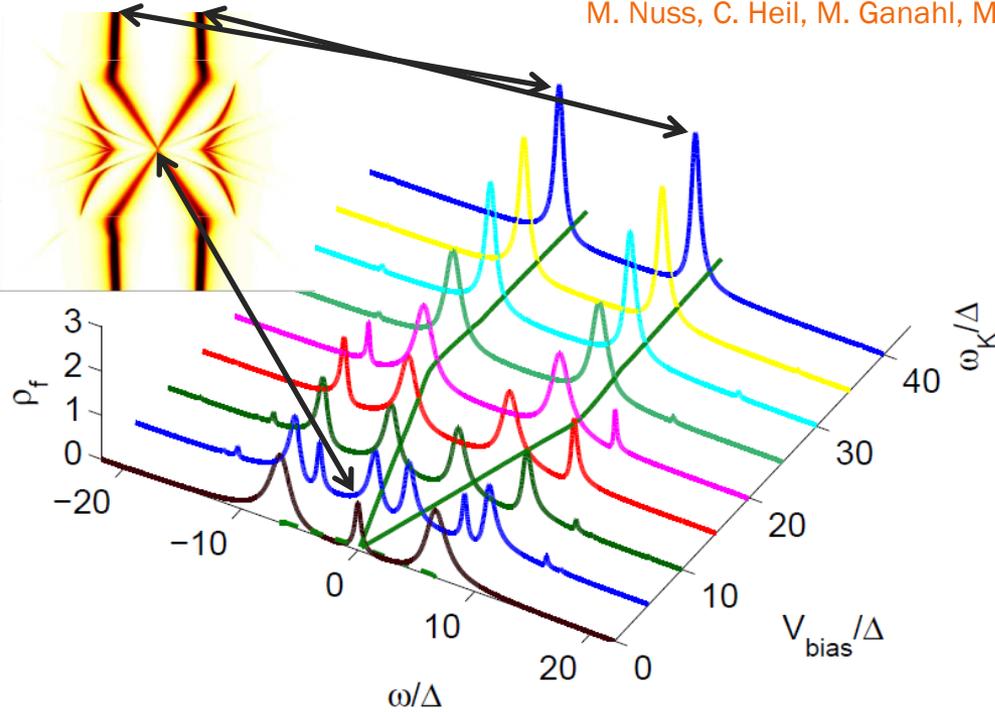
A: splitting of Kondo resonance: *linear + U dep.*

Equilibrium
 Kondo resonance @ $\omega = 0$
 inc. excitations @ $\omega \approx \pm \frac{U}{2}$

B: finite size eff. in Hubbard band
C: lead-band effects

Kondo resonance under bias

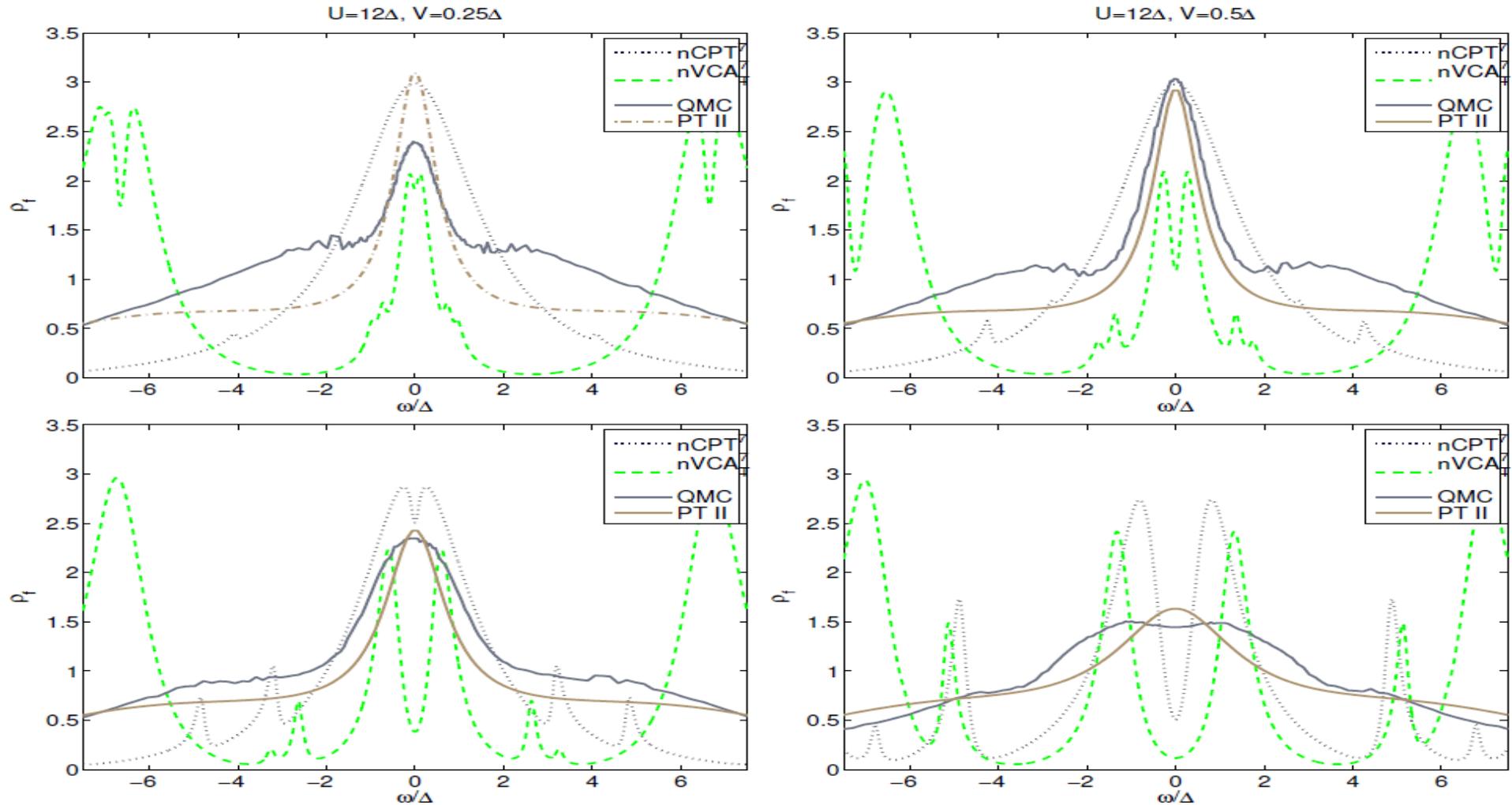
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VCA: immediate linear splitting

pinning not exactly at chemical potential but U dependent

Kondo resonance under bias

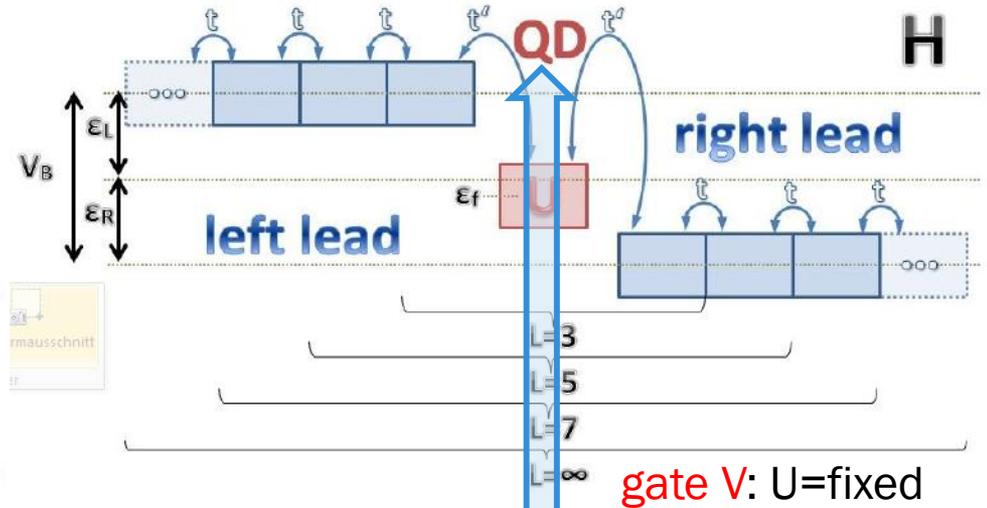
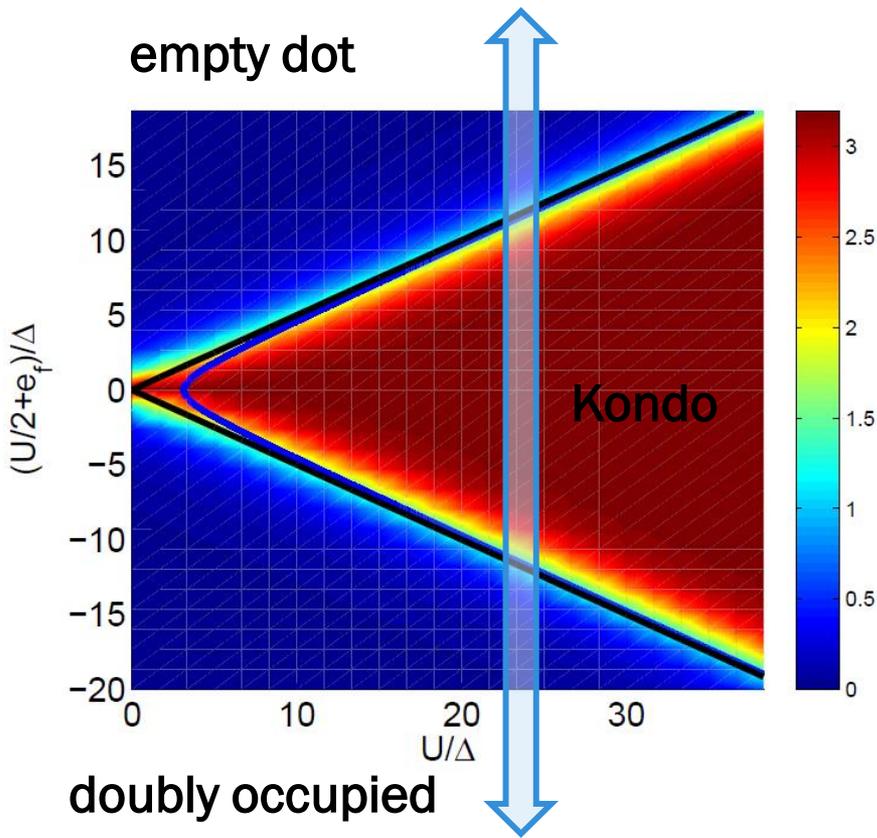


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A. Dirks, J. E. Han, M. Jarrell, T. Pruschke (2012)

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applying a gate voltage

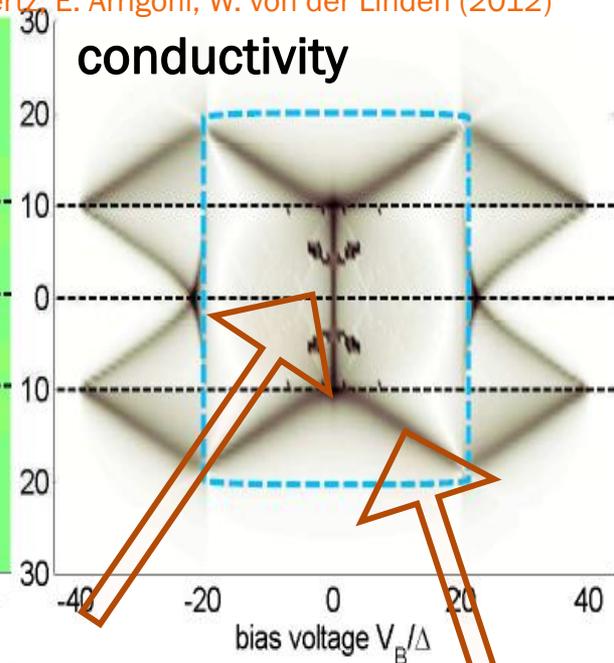
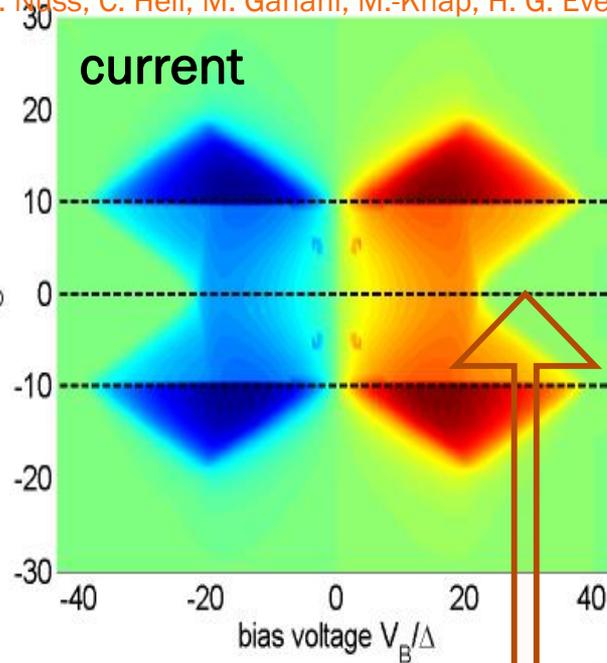
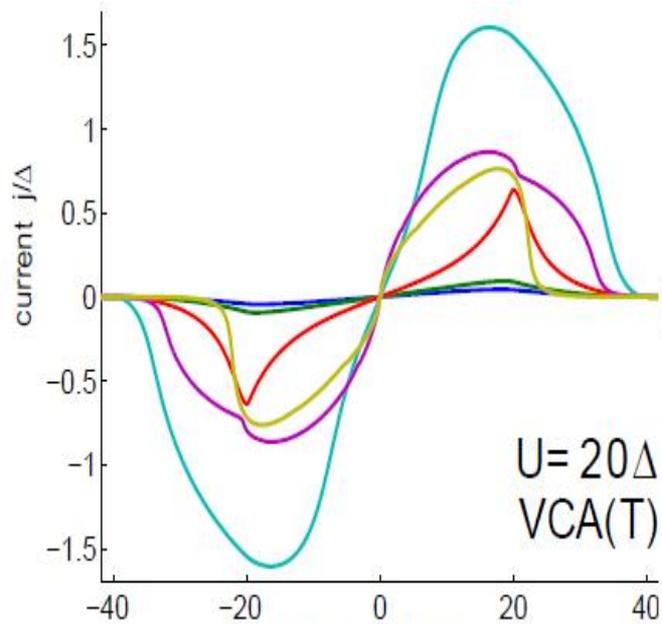


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applying a gate voltage

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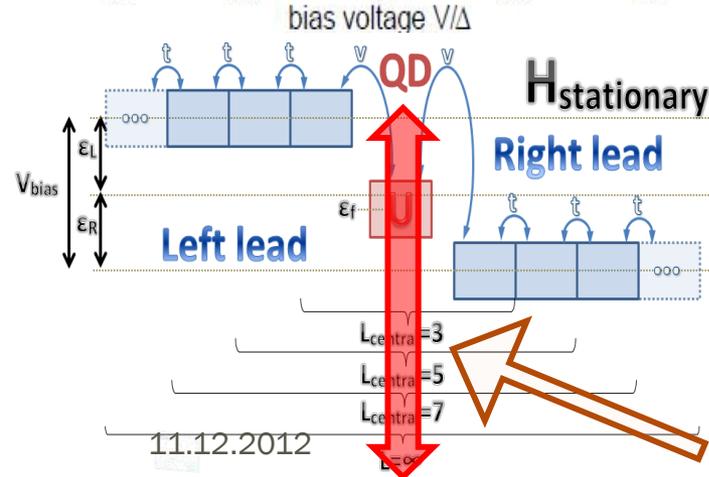


Kondo

coulomb blockade

results from previous slides $V_{gate}=0$

gate V : U =fixed



Conclusions + Outlook

∞ steady-state: Quantum Dot

- good current density up to intermediate U
- agrees with TEBD benchmark
- linear U dep. splitting of Kondo resonance
- Kondo regime + Coulomb blockade
- $nVCA \gg nCPT$: variational feedback crucial

∞ non-equilibrium Variational Cluster Approach

- applicable to any fermionic bosonic lattice Hamiltonian
- benchmark on SIAM good
- more complex models, interactions, versatile and flexible
- dynamic quantities out of equilibrium (real)
- realistic materials: combine with ab-initio

Thank you!

maybe some day:



martin.nuss@student.tugraz.at

itp^{cp}

TUG

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We gratefully acknowledge support from the Austrian Science Fund (FWF) P24081-N16 as well as the Vienna Scientific Cluster (VSC).

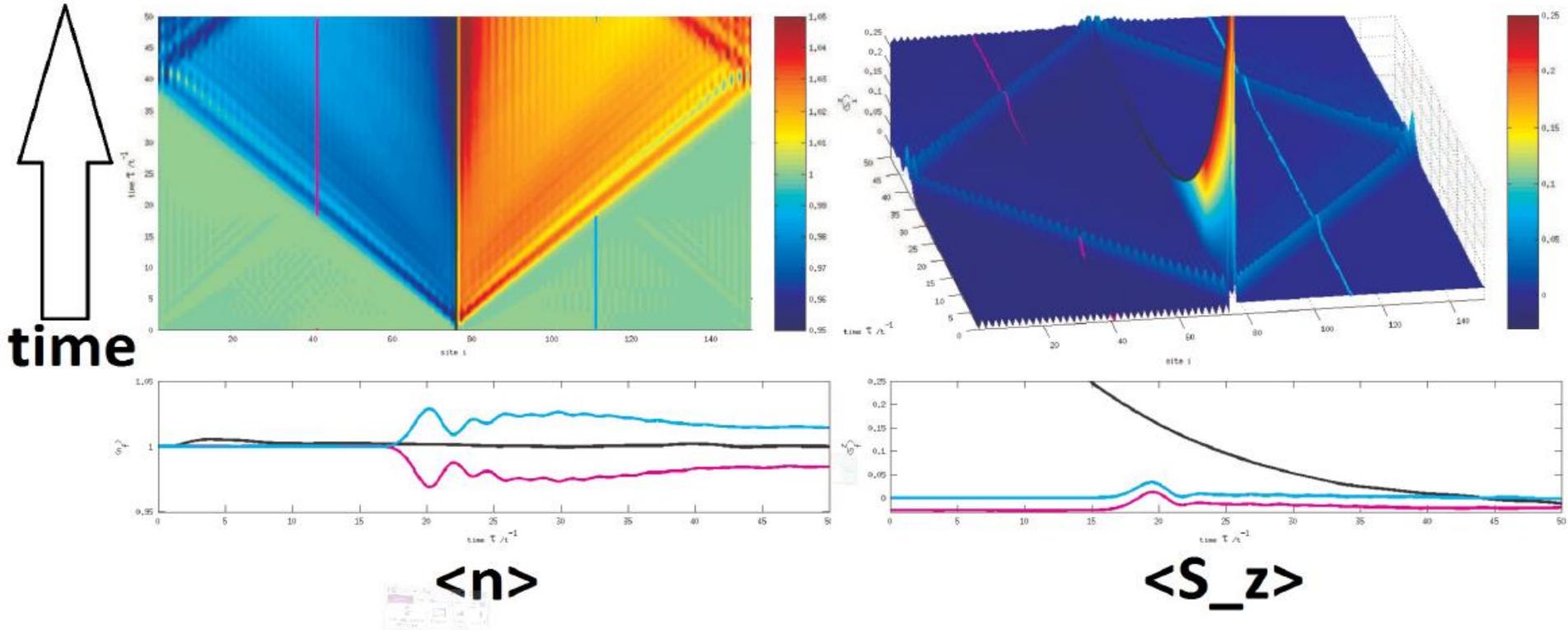
Time evolution of a quantum dot

∞ 4 ∞

particle number and spin projection

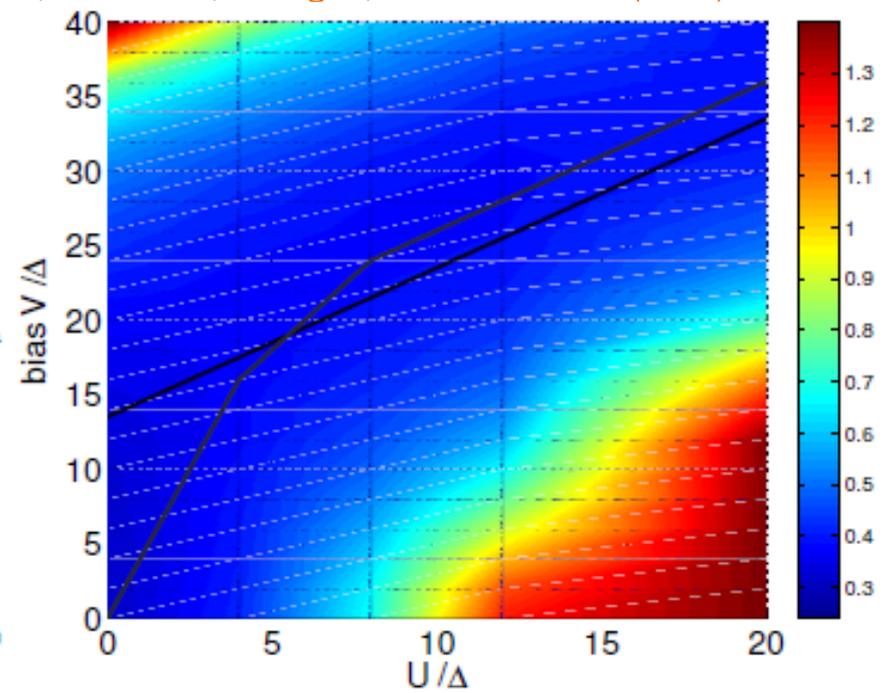
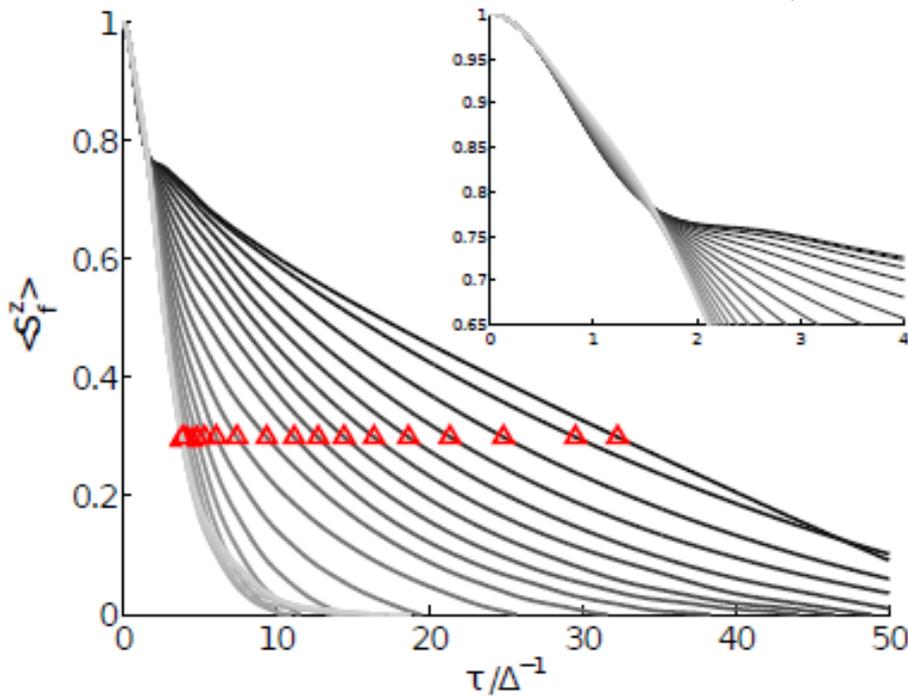
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$$L = 150, U = 20\Delta, V_{\text{bias}} = 14\Delta$$



Spin relaxation

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Entanglement

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