## Comment on "Quantum Monte Carlo Evidence for Superconductivity in the Three-Band Hubbard Model in Two Dimensions"

In a recent Letter, Kuroki and Aoki [1] presented quantum Monte Carlo (QMC) results for pairing correlations in the three-band Hubbard model, which describes the Cu- $d_{x^2-y^2}$  and O- $p_{x,y}$  orbitals present in the CuO<sub>2</sub> planes of high- $T_c$  materials. In this Comment we concentrate on the parameter set  $U_d = 3.2t_{pd}$ ,  $\Delta = 2.7t_{pd}$ ,  $t_{pp} = -0.4t_{pd}$ . For this parameter choice, Kuroki and Aoki see a maximal increase in the  $d_{x^2-y^2}$  pairing correlations which they associate with a signature of offdiagonal long-range order (ODLRO). We argue that:

(i) The above parameter set is not appropriate for the description of high- $T_c$  materials since it does not satisfy the minimal requirement of a charge-transfer gap at half-filling. To illustrate this point, we have calculated with QMC methods the average hole number as a function of the chemical potential:  $\langle n \rangle (\mu)$ . Our results, which are plotted in Fig. 1, show a vanishingly small charge-transfer gap (i.e.,  $\Delta_{ct} < 0.07t_{pd}$ ). In contrast, for a physical parameter set [2], one obtains a sizable charge-transfer gap which is detectable from the plateau in the  $\langle n \rangle (\mu)$  curve (see inset, Fig. 1). For the latter parameter set, a number of normal state properties were shown to successfully reproduce experimental data [3]. However, despite intensive numerical efforts, no ODLRO was unambiguously detected [4].

(ii) The observed increase in the  $d_{x^2-y^2}$  channel (Fig. 2 in Ref. [1]) is dominantly produced by the pair-field correlations without the vertex part [5]. To prove this point we have calculated the pair-field correlations in the *d*-wave channel summed over distances **r** with  $|r_x|, |r_y| < R$  [ $S_d(R)$ ] [see Fig. 2(a)]. As in Ref. [1], an increase as a function of *R* can be seen. However, the vertex contribution to the pair-field correlations, which is the



FIG. 1.  $\langle n \rangle(\mu)$  on a 4 × 4 lattice at  $\beta t_{pd} = 10,15$  for the parameter set  $U_d = 3.2t_{pd}, \Delta = 2.7t_{pd}, t_{pp} = -0.4t_{pd}$ . Inset:  $\langle n \rangle(\mu)$  on a 4 × 4 lattice at  $\beta t_{pd} = 10$  for  $U_d = 6t_{pd}, \Delta = 4t_{pd}, t_{pp} = 0.0$  [2].



FIG. 2. (a)  $S_d(R)$  on an  $8 \times 8$  lattice at zero temperature and  $U_4 = 3.2t_{pd}, \Delta = 2.7t_{pd}, t_{pp} = -0.4t_{pd}$ . (b) Vertex contribution to  $S_d(R)$  shown in (a).

relevant quantity, is an order of magnitude smaller and shows—within our numerical accuray—no significant increase as a function of R [see Fig. 2(b)]. Hence the claim of evidence of ODLRO is not justified.

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- [1] K. Kuroki and H. Aoki, Phys. Rev. Lett. **76**, 4400 (1996).
- [2] G. Dopf, A. Muramatsu, and W. Hanke, Phys. Rev. Lett. 68, 353 (1992).
- [3] G. Dopf, J. Wagner, P. Dieterich, A. Muramatsu, and W. Hanke, Phys. Rev. Lett. 68, 2082 (1992).
- [4] F. F. Assaad, W. Hanke, and D. J. Scalapino, Phys. Rev. B 49, 4327 (1994).
- [5] S. R. White, D. J. Scalapino, R. L. Sugar, N. E. Bickers, and R. T. Scalettar, Phys. Rev. B **39**, 839 (1989).