

Elektrodynamics Exercises WS 2015-16

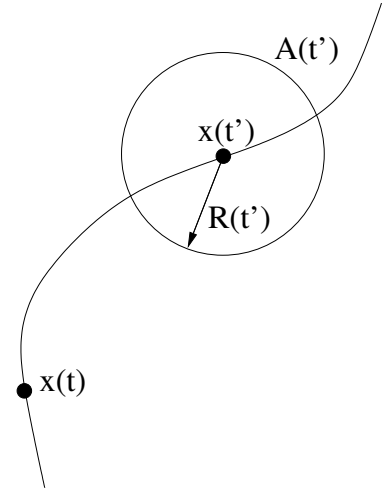
Problem set 4, 04.12.2015

Problem 1: Radiation of a charged particle

Eq. (15.72) in the lecture notes defines the Poynting vector for the radiation field:

$$\mathbf{S}_s = \left[\frac{q^2 \mathbf{N}}{16\pi^2 \epsilon_0 c^3 \kappa^6 R^2} (\mathbf{N} \times [(\mathbf{N} - \mathbf{v}/c) \times \mathbf{a}])^2 \right]_{ret}$$

where $\mathbf{N} = \mathbf{R}/R$. The energy that flows per time unit through the surface element $d\mathbf{A}(t')$ is given by $\mathbf{S}_s(t') \cdot d\mathbf{A}(t')dt$. Because all properties are expressed in the retarded time t' , and not in t , we have to transform all quantities from dt to dt' . This gives the flow of energy as $dP(t')dt' = \mathbf{S}_s(t') \cdot d\mathbf{A}(t')(dt/dt')dt'$.



- a) Calculate the factor (dt/dt') and show that

$$\frac{dP}{d\Omega} = \left[\frac{\mu_0 q^2}{(4\pi)^2 c} \cdot \frac{|\mathbf{N} \times [(\mathbf{N} - \mathbf{v}/c) \times \mathbf{a}]|^2}{\kappa^5} \right]_{ret}$$

with $d\Omega(t')$ the element of the solid angle, $\kappa = 1 - \mathbf{N} \cdot \mathbf{v}/c$. The retarded time t' is determined by the condition $|\mathbf{R}(t')| = c(t - t')$.

Consider now the following cases:

- b) Slow Motion. Let the θ be the angle between the acceleration \mathbf{a} and \mathbf{R} . Consider now small velocities $v \ll c$, $\kappa \approx 1$, and

b1) show that $\frac{dP}{d\Omega} = \left[\frac{\mu_0 q^2}{(4\pi)^2 c} |\mathbf{a}|^2 \sin^2 \theta \right]_{ret}$.

b2) integrate $\frac{dP}{d\Omega}$ over the solid angle. Show that $P = \left[\frac{\mu_0 q^2}{6\pi c} |\mathbf{a}|^2 \right]_{ret}$.

- c) Linear Motion Show that the power depends in the following way on the angle:

$$\frac{dP}{d\Omega} = \left[\frac{\mu_0 q^2}{(4\pi)^2 c} |\mathbf{a}|^2 \frac{\sin^2 \theta}{(1 - v \cdot \cos \theta/c)^5} \right]_{ret}$$

where θ is again angle between \mathbf{a} and \mathbf{R} . Take into account that \mathbf{v} and \mathbf{a} are parallel.

- c1) Show that for $v \ll c$ you get back equation (b1).

- c2) Show that the maximum of $dP/d\Omega$ is reached for $\cos\theta_{max} = \frac{\sqrt{1+15(v/c)^2}-1}{3(v/c)}$.
Hint: Substitute $\cos\theta \equiv x$ and find maximum with respect to x .
- c3) Do an expansion of above equation up to second order in $\beta = v/c \ll 1$ and show $\theta_{max} = \frac{\pi}{2} - \frac{5}{2}\beta$.
Hint: Please use Taylor expansion of the cosine at $\theta = \pi/2$.
- c4) Expand now $\cos\theta_{max}$ in terms of γ^{-1} up to γ^{-2} for the case $v \approx c$, $\gamma^{-2} = 1 - (v/c)^2 \ll 1$. Show that $\theta_{max} = \frac{1}{2}\gamma^{-1}$. What does this result mean for high (relativistic) velocities, in what direction does the particle radiate?
Hint: Please use Taylor expansion of the cosine at $\theta = 0$.
- c5) Calculate the power P by integrating over the solid angle $d\Omega$ and show that $P = \left[\frac{\mu_0 q^2}{6\pi c} |\mathbf{a}|^2 \gamma^6 \right]_{ret}$.
- d) Circular motion. Show that the angular dependence of the radiated power is

$$\frac{dP}{d\Omega} = \left[\frac{\mu_0 q^2}{(4\pi)^2 c} \frac{|\mathbf{a}|^2}{(1 - v \cdot \cos\theta/c)^3} \left[1 - \frac{\sin^2\theta \cos^2\phi}{\gamma^2(1 - v \cdot \cos\theta/c)^2} \right] \right]_{ret}$$

Here θ and ϕ are polar and azimuthal angles respect to \mathbf{v} . Find the values of θ, ϕ , for which $dP/d\Omega$ reaches its maximum. Show that this radiation is in forward direction, i.e. in the direction of the velocity \mathbf{v} of the particle (synchrotron radiation). Take into account that \mathbf{v} and \mathbf{a} are orthogonal vectors.

Problem 2: Dipole antenna

We consider an antenna with the following current density (I_0 is a constant):

$$\mathbf{j} = I_0 \operatorname{Re} \delta(x)\delta(y)\theta(L/2 - |z|)\mathbf{e}_z e^{-i\omega t}.$$

- a) Calculate the vector potential $\mathbf{A}(\mathbf{r}, t)$ (in Lorentz gauge) to leading order in $1/r$. Note that L/λ ($\lambda = 2\pi c/\omega$) is not small!
Hint: You can substitute $|\mathbf{r} - \mathbf{r}'|$ in the denominator by $|\mathbf{r}|$. In the exponent, however, you have to use $|\mathbf{r} - \mathbf{r}'| \approx |\mathbf{r}| - \mathbf{r}' \cdot \mathbf{r}/r$.
- b) Calculate the field $\mathbf{B}(\mathbf{r}, t)$, the Poynting vector and its time average $\bar{\mathbf{S}}$.
Hint: for large \mathbf{r} , $\nabla \rightarrow ik\mathbf{N}$, with $\mathbf{N} = (\mathbf{r} - \mathbf{r}')/|\mathbf{r} - \mathbf{r}'|$.
- c) Determine the angular dependence of $\bar{\mathbf{S}}$ for a $\lambda/2$ antenna, i.e. with $L = \lambda/2$. Show furthermore that for short antennas, i.e. $L \ll \lambda$, the angular dependence is given by the characteristic of a Hertzian dipole (Eq. 16.24 in the lecture notes):

$$\bar{\mathbf{S}} = \frac{\mu_0}{16\pi^2 c} \omega^4 d^2 \frac{\sin^2\theta}{2r^2} \mathbf{N}.$$

- d) Calculate for a long antenna, $L \gg \lambda$, the total power of radiation to infinity.